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SOURCE

## Spin dynamics in the hyperkagome compound $\text{Gd}_3\text{Ga}_5\text{O}_{12}$

Pascale Deen

O. Petrenko<sup>1</sup>, H. Mutka<sup>2</sup>, T. Fennell<sup>3</sup>, G. Balakrishnan<sup>1</sup>,  
B. Rainford<sup>3</sup>, C. Ritter<sup>2</sup>, L. Capogna<sup>2</sup>, B. Canals<sup>4</sup>, M. Zhitomirsky<sup>5</sup>

<sup>1</sup> University of Warwick, UK.

<sup>2</sup> Institut Laue Langevin, Grenoble, France.

<sup>3</sup> University of Southampton, UK.

<sup>4</sup> CNRS - Institut Néel, Grenoble, France.

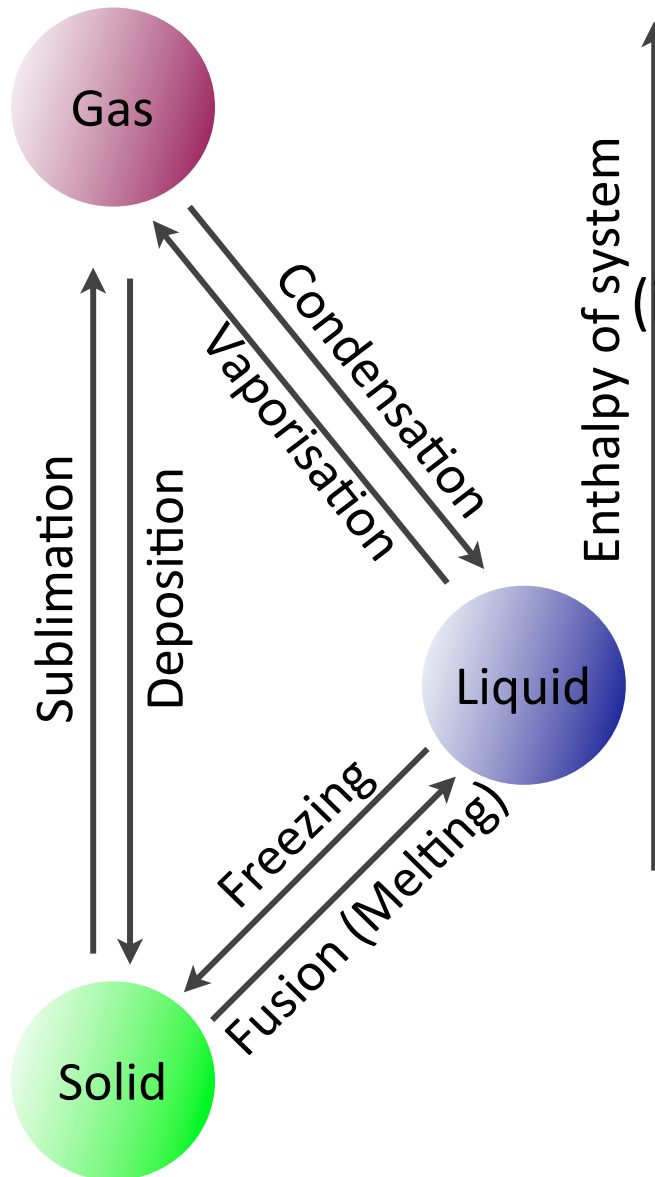
<sup>5</sup> CEA- Grenoble France.



## Outline

- Phase transitions, frustration and dynamics
- Pyrochlores versus Kagome compounds
- 3D Kagome  $\text{Gd}_3\text{Ga}_5\text{O}_{12}$  :
  - a background
  - Inelastic neutron scattering
- What does it mean.

## Phase transitions



### Universality

Liquids = Magnetism

$$(T - T_c)^\alpha$$

Liquid-gas critical point (independent of composition) FM phase transitions uniaxial magnets

### Same universality class

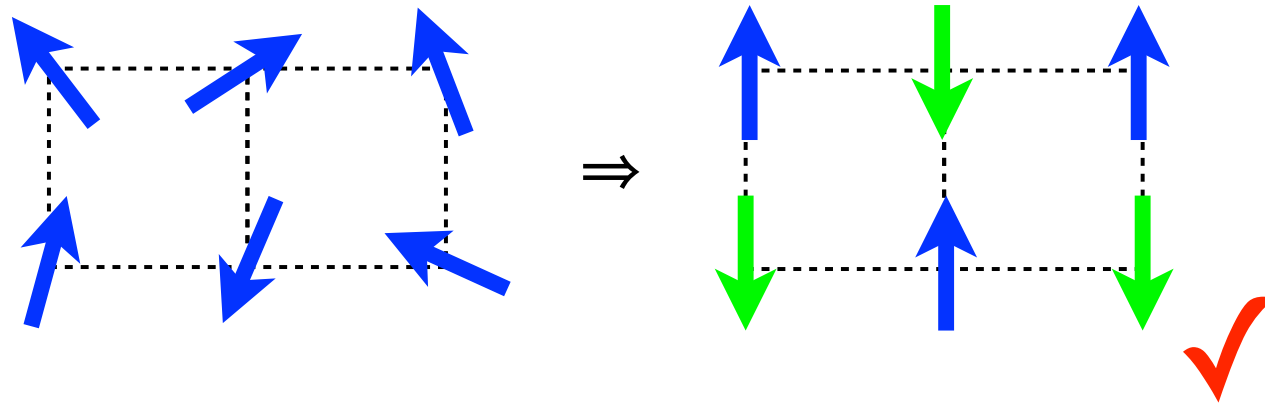
*Universal thermodynamic properties of a system near a phase transition depend only on a small number of features, such as dimensionality and symmetry, and are insensitive to the underlying microscopic properties of the system.*



## Magnetic phase transitions

As  $T \Rightarrow 0$  , Divergent correlation lengths

Phase transition to long range order



Minimise all interactions to achieve a ground state

$$\sigma_1 + \sigma_2 + \sigma_3 + \dots = 0$$





# Frustration $\Rightarrow$ zero point entropy $\Rightarrow$ degeneracy

## Antiferromagnetism. The Triangular Ising Net

G. H. WANNIER

*Bell Telephone Laboratories, Murray Hill, New Jersey*

(Received February 11, 1950)

In this paper the statistical mechanics of a two-dimensionally infinite set of Ising spins is worked out for the case in which they form either a triangular or a honeycomb arrangement. Results for the honeycomb and the ferromagnetic triangular net differ little from the published ones for the square net (Curie point with logarithmically infinite specific heat). The triangular net with antiferromagnetic interaction is a sample case of antiferromagnetism in a non-fitting lattice. The binding energy comes out to be only one-third of what it is in the ferromagnetic case. The entropy at absolute zero is finite; it equals

$$S(0) = R \frac{2}{\pi} \int_0^{\pi/3} \ln(2 \cos \omega) d\omega = 0.3383R.$$

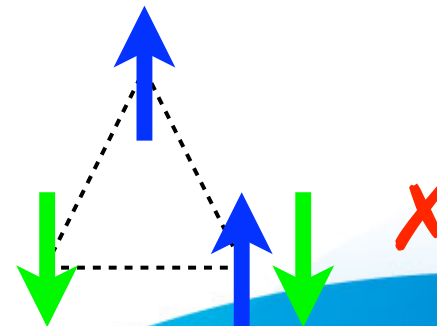
The system is disordered at all temperatures and possesses no Curie point.

### Geometric Frustration

AF interactions  $J > 0$

Ising spins

$(\sigma_i = \pm 1)$



Entropy = 0.34 kb /spin at T= 0 K.

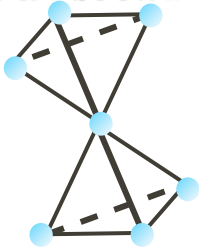


# Frustration $\Rightarrow$ zero point entropy $\Rightarrow$ degeneracy

PHYSICAL REVIEW

VOLUME 102, NUMBER 4

MAY 15, 1956

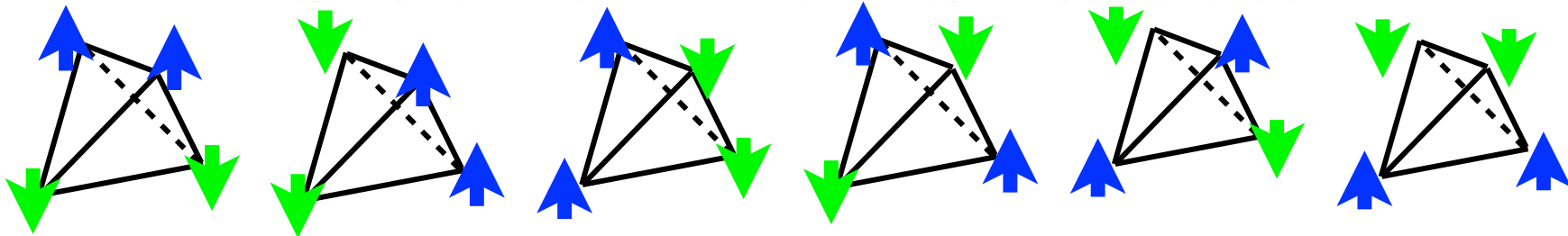


## Ordering and Antiferromagnetism in Ferrites

P. W. ANDERSON

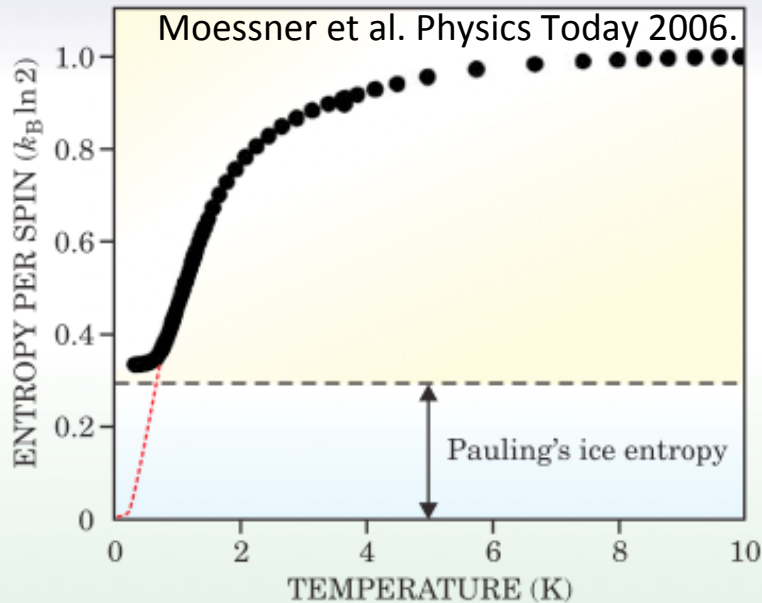
*Bell Telephone Laboratories, Murray Hill, New Jersey*

(Received January 9, 1956)



A. P. Ramirez et al., *Nature* 399, 333 (1999).

Moessner et al. *Physics Today* 2006.



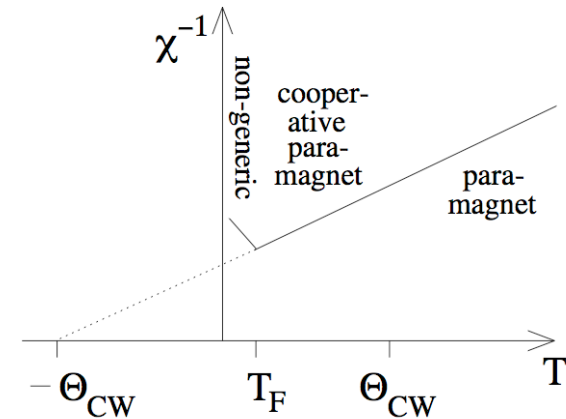
*Nearest neighbour exchange cannot lead to long range order but there is short range order with finite zero point entropy*  
 Manifold of connected ground states

# Frustration $\Rightarrow$ zero point entropy $\Rightarrow$ degeneracy

## Signatures

- No order down to lowest temperatures  
 $T \ll J(\text{Near neighbour}) \sim \theta_{CW}$
- Zero point entropy. Degenerate ground state

'Susceptibility fingerprint'  
of frustration

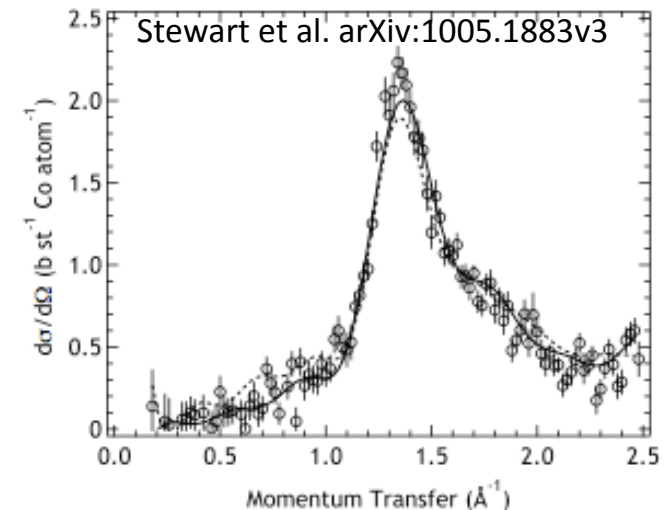


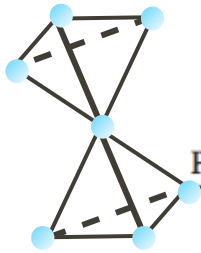
- Short range correlations in space

Neutron scattering

$$\Delta Q = 1/\xi$$

Dynamic behaviour?





## Low-temperature properties of classical geometrically frustrated antiferromagnets

R. Moessner and J. T. Chalker

PRL 102, 237206 (2009)

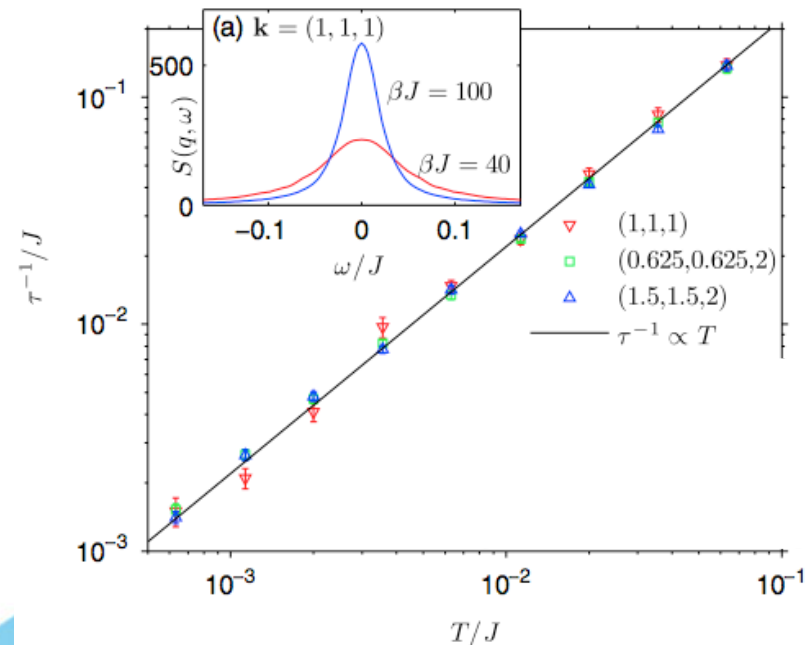
PHYSICAL REVIEW LETTERS

week ending  
12 JUNE 2009

## Spin Dynamics in Pyrochlore Heisenberg Antiferromagnets

P. H. Conlon\* and J. T. Chalker

- Time dependence of spin-spin autocorrelation function.  
 $\langle \mathbf{S}_R(0) \mathbf{S}_R(t) \rangle = \exp(-cTt),$
- FT = Lorentzian scattering in frequency space with  $\Delta E \sim T$ .  
Depends on T not J.
- Spin correlations relax at a rate independent of the wave vector and proportional to the temperature.

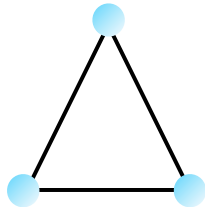




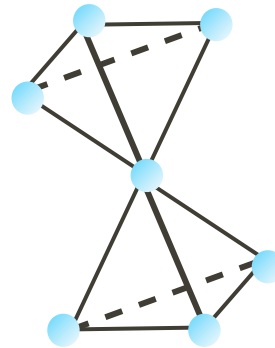
Frustration  $\Rightarrow$  zero point entropy  $\Rightarrow$  degeneracy

$$S > 1/2 = \text{not quantum}$$

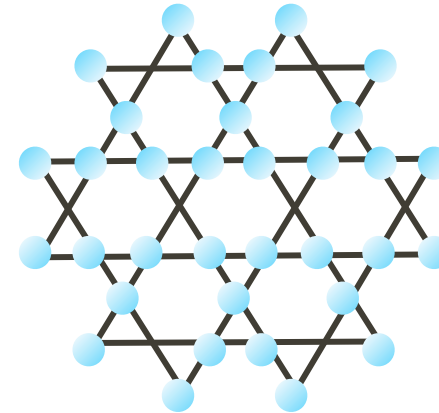
Triangle



Pyrochlore



Kagome



Spin liquid, spin ice    **How does it compare?**

Magnetic analogue of ice and liquid.

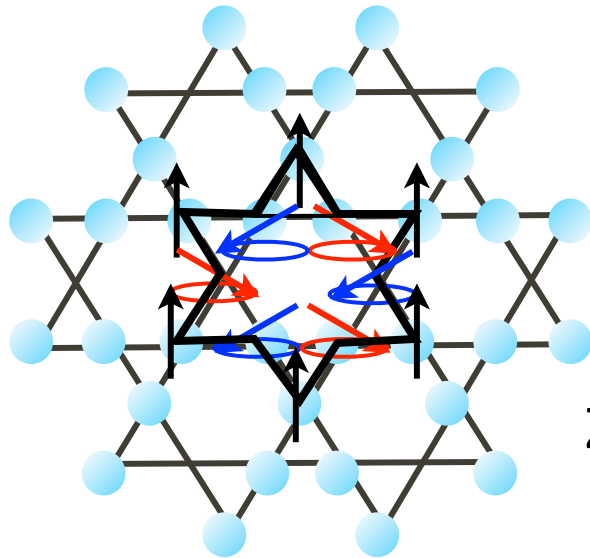
Dynamically: Dynamic gapless excitations (degeneracy) to 0 K.

Universal dynamic behaviour?

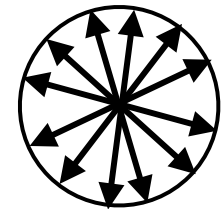


## Kagome lattice

Prototype of geometrical frustration  
Juxtaposition of theory and experiment



NN Heisenberg AF : Infinite degenerate number of spin configurations



Constraint of AF Spins on a Kagome lattice constraint =  $120^\circ$

Zero energy modes (3 spin, 6 spin, xx spin)

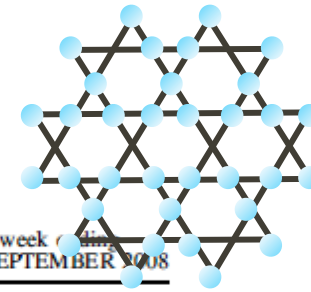
Non-dispersive and costs no energy.

Rotations of finite number of spins:

No correlations in space.

Hall mark of Kagome dynamics.





## Propagation and Ghosts in the Classical Kagome Antiferromagnet

J. Robert,<sup>1</sup> B. Canals,<sup>2</sup> V. Simonet,<sup>2</sup> and R. Ballou<sup>2</sup>

### Nearest neighbour interactions only

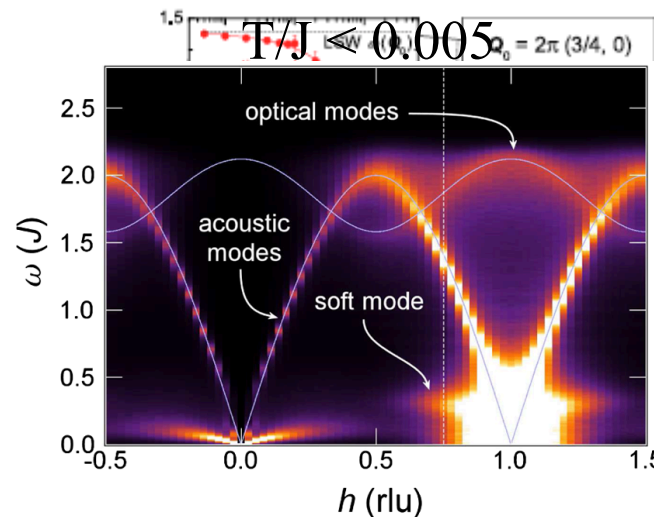


FIG. 2 (color online). Temperature weighted scattering function vs energy at different temperatures for  $\mathbf{Q}_0 = 2\pi(3/4, 0)$ . Inset: position of soft and acoustic modes vs temperature; error bars obtained from several fitting processes.

$T/J_{NN}$	Expectation
$> 0.2$	QE signal only
$0.2-0.01$	Broad excitation
$< 0.01$	Splits into acoustic and soft mode
$< 0.005$	+ Optical mode

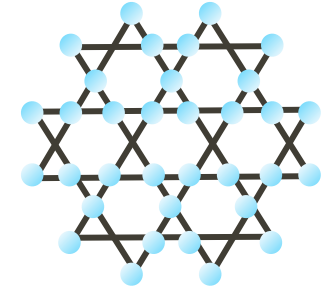
Sufficient temporal and spatial stiffness gives rise to soft modes. They fall to zero energy as  $T \rightarrow 0$  K.



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## Dynamics to date

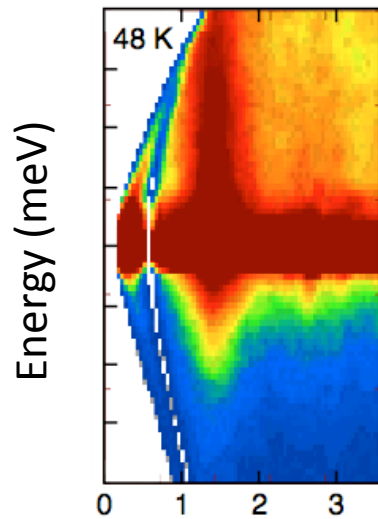
Classical Heisenberg AF on a Kagome lattice  
Gapless excitations extend far out in energy



Spatial dependence: SRO

$Y_{0.5}Ca_{0.5}BaCo_4O_7$ :

J. R. Stewart et al. arXiv:1005.1883v3



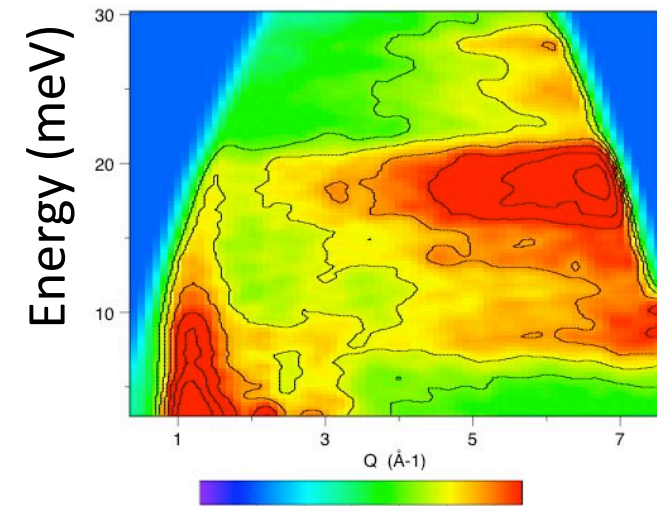
Wavevector transfer ( $\text{\AA}^{-1}$ )

- Remains fluctuation down to lowest temperatures.
- No gap in the excitations - no internal energy barrier

Deuterium Jarosite :

B. Fåk, *et al.*

Europhysics Letters (2008) 81, 17006

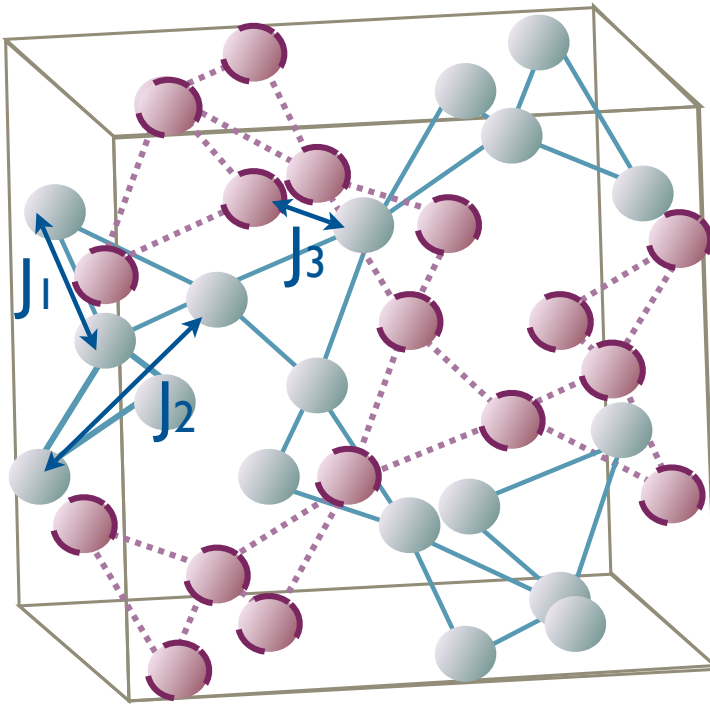


Wavevector transfer ( $\text{\AA}^{-1}$ )





## Spin Dynamics in the Hyperkagome (3D Kagome) $\text{Gd}_3\text{Ga}_5\text{O}_{12}$



- Each RE spin has 4 NN
- NN form lattice of equilateral triangles.
- 2 Interpenetrating corner sublattices, spin do NOT lie in same plane.
- No inversion symmetry.

### Bulk measurements

Onn et al. Phys.Rev. (1966), 156, 663.

- $J_1$  = NN exchange energy  $\sim 1.68\text{K}$ ,  $J_2$  &  $J_3 = J_1/1000$
- Dipole energy  $\sim 0.7\text{ K}$  - anisotropic
- Curie Weiss  $\sim -2.8\text{K}$
- No long range order down to 25 mK.

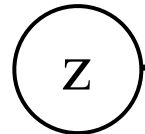
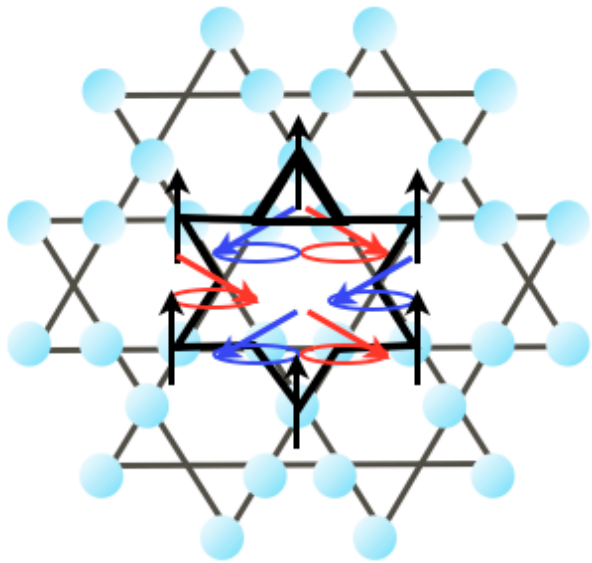


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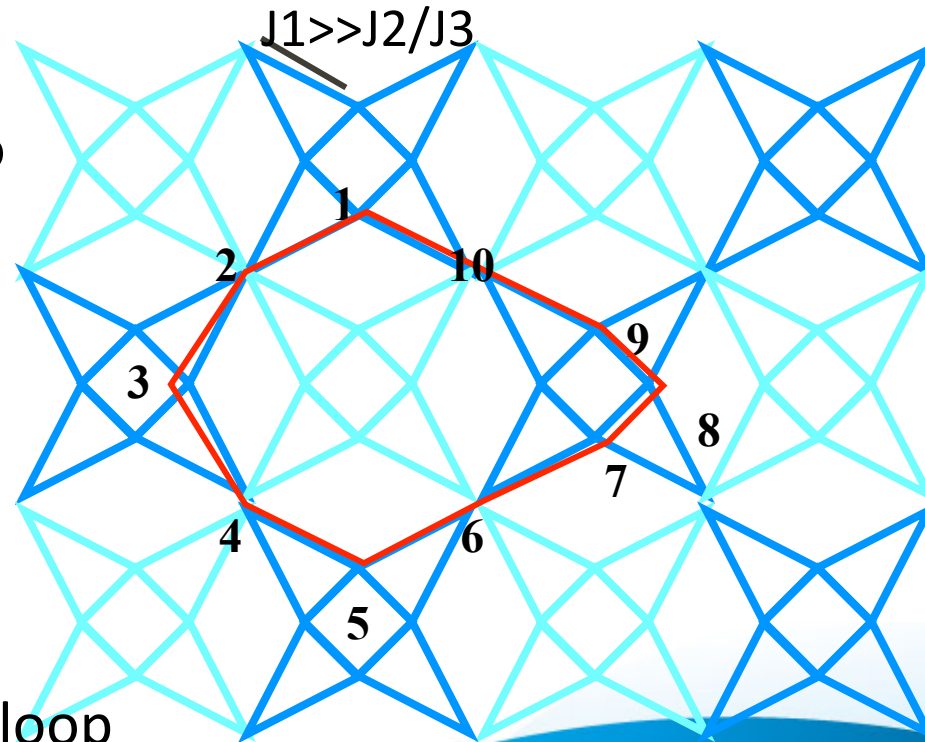


Soft modes

Kagome



Hyper Kagome



3 spin loop  
10 spin loops  
xx spin loop



## Gd<sub>3</sub>Ga<sub>5</sub>O<sub>12</sub>

Schiffer et al. Phys. Rev. Lett (1994),73, 2500

Schiffer et al. Phys. Rev. Lett. (1995), 74, 22380

Kinney et al. J. App. Physics (1979), 50, 2115.

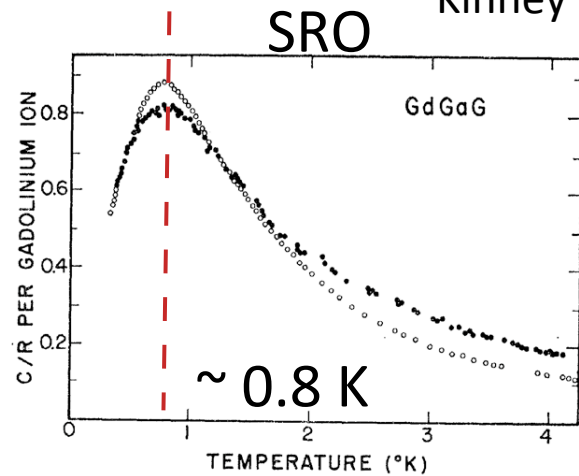


FIG. 5. Specific heat of two samples of GdGaG. ● Sintered sample, ○ cluster of single crystals.

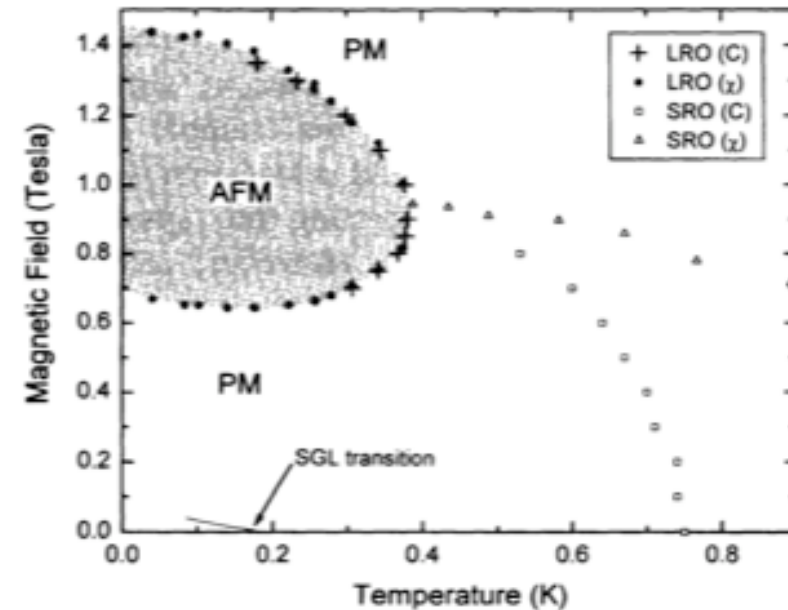
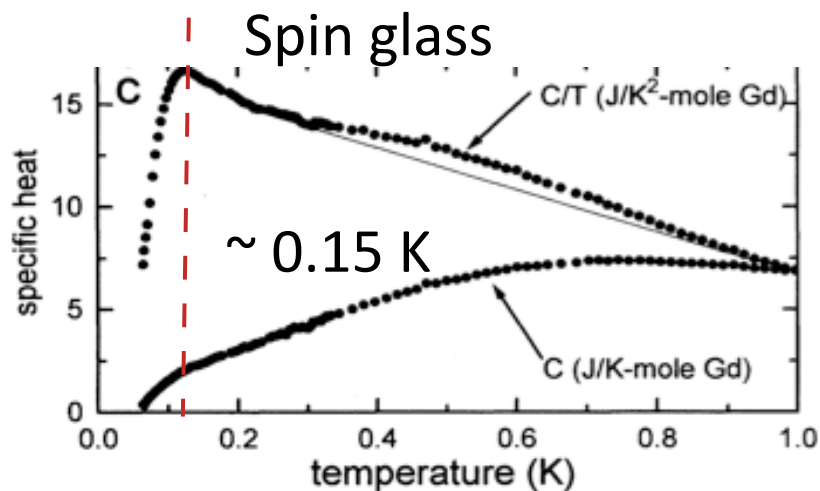


FIG. 2. The magnetic phase diagram of GGG. The long range order peaks which define the antiferromagnetic (AFM) phase boundary and the short range order peaks in the paramagnetic (PM) phase are shown. The line of maxima in  $\chi$  vs  $T$  corresponding to the spin-glass-like freezing transition is also shown.



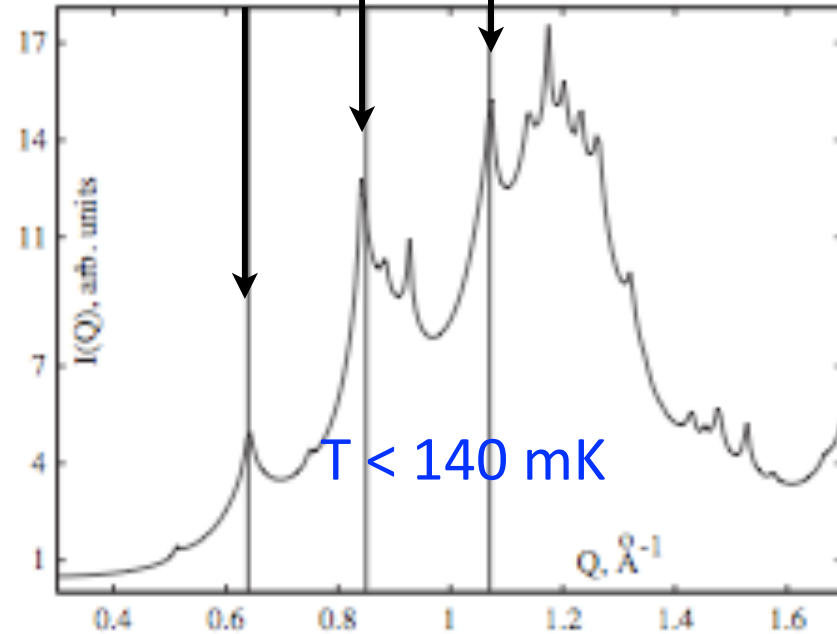
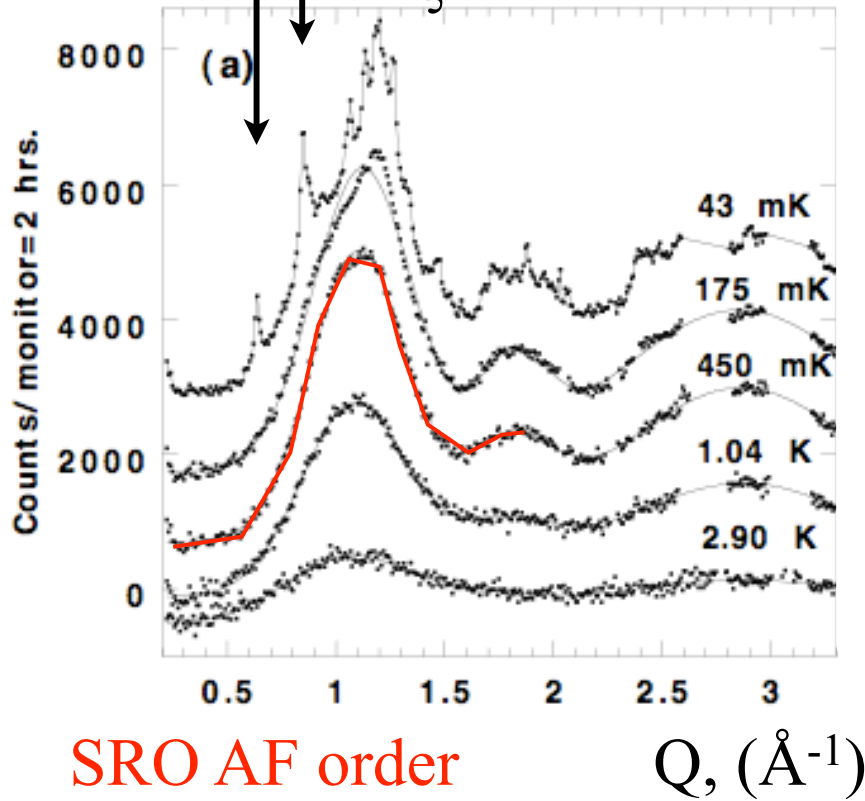
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# Neutron scattering

## Temperature dependence

$T < 140 \text{ mK}$   
 $\xi \sim 100 \text{ \AA}$

need J1, J2, J3 and D



*We show that it is crucial to treat accurately the long-range nature of the magnetic dipolar interactions to allow for a determination of the small exchange energy scales involved in the selection of the experimental ordering wave vector.*

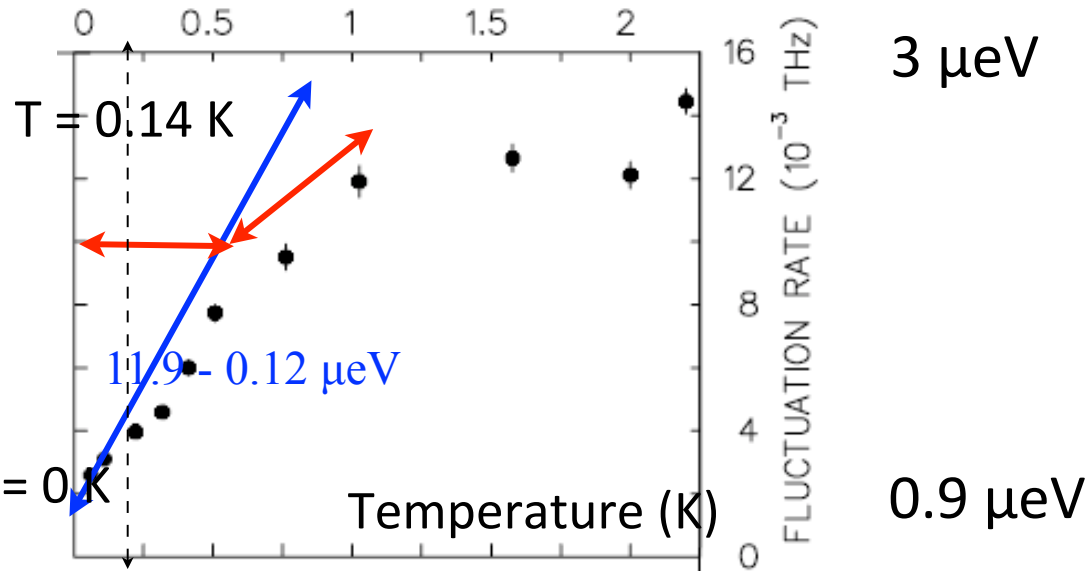
Petrenko et al. Phys. Rev. Lett. (80), 4570

T. Yavors'kii et al. Phys. Rev. Lett (97), 267203, 2006



## Dynamic behaviour - spin relaxation

### Pyrochlore or Kagome?



Finite dynamics at  $T = 0$  K

$\mu\text{SR}$ : Absence of LRO (indirect study)

Dunsiger et al. Phys. Rev. Lett. 85, 3504 (2004)

Marshall J. Phys. Cond Matt. 14, L157 (2002)

Mossbauer: Linear decrease (11.9 - 0.12  $\mu\text{eV}$ ) (direct study) planar confinement  
Bonville Phys. Rev. Lett. 92, 167202 (2004)

Magnetic hole burning: S. Gosh Phys. Rev. Lett. 101, 157205 (2008)

$T < 0.14$  K protected spin clusters (10 ring)

around impurity centers - pico eV dynamics !

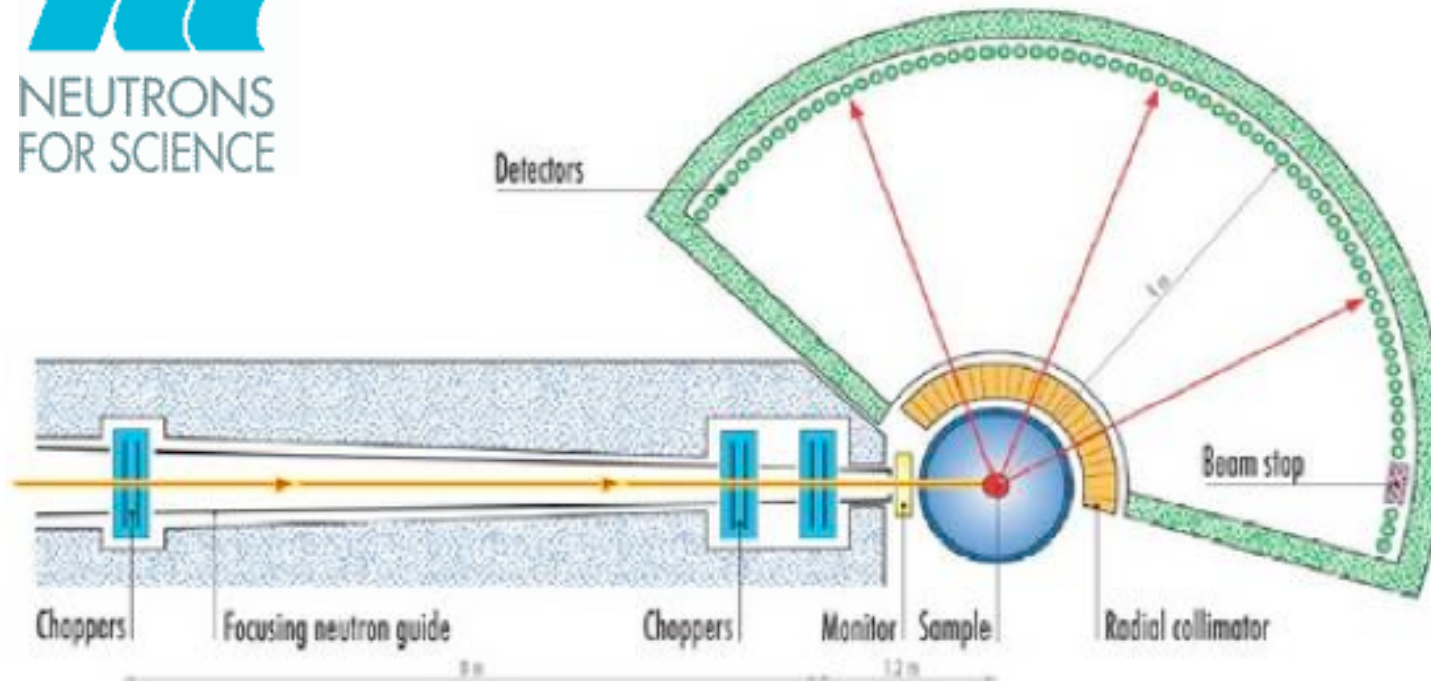




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# IN5 - Cold neutron TOF spectrometer

## Gain Q and E dependence



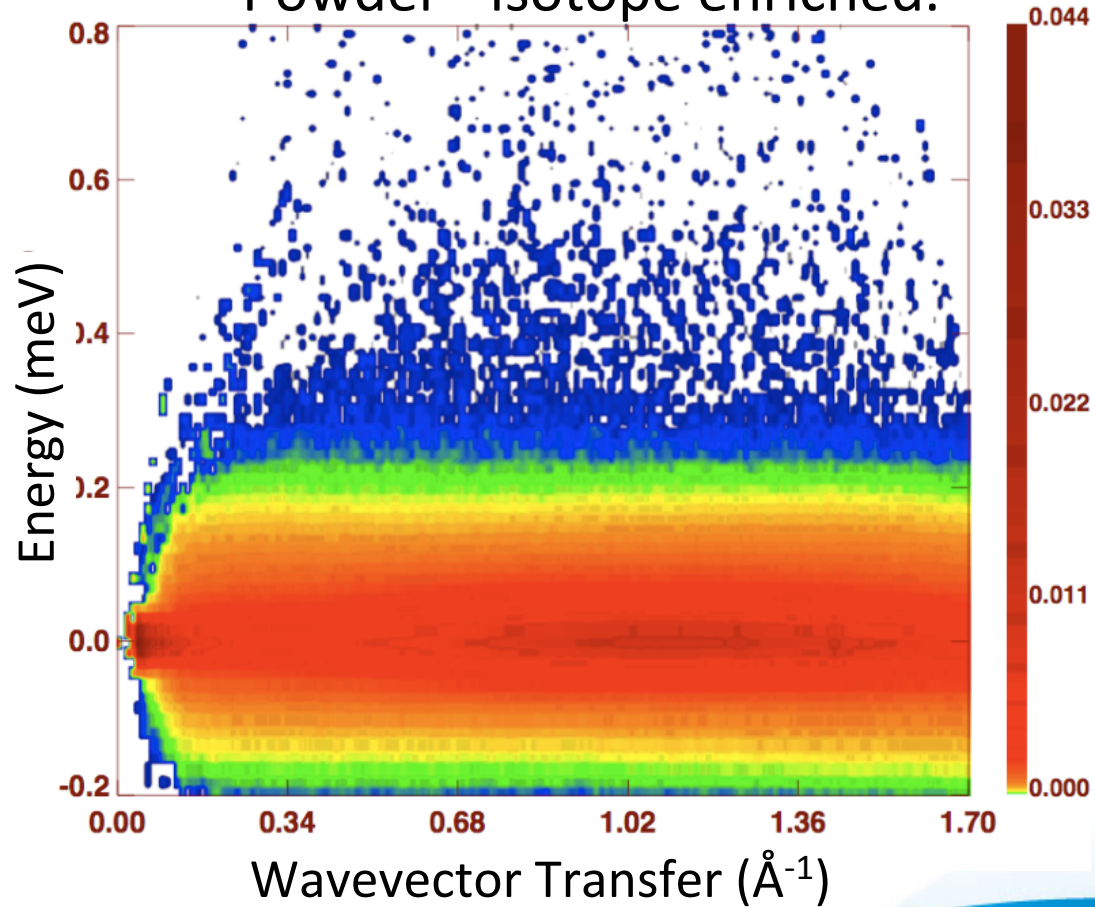


# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

2 K

Powder - Isotope enriched.

FT of spin spin autocorrelation.

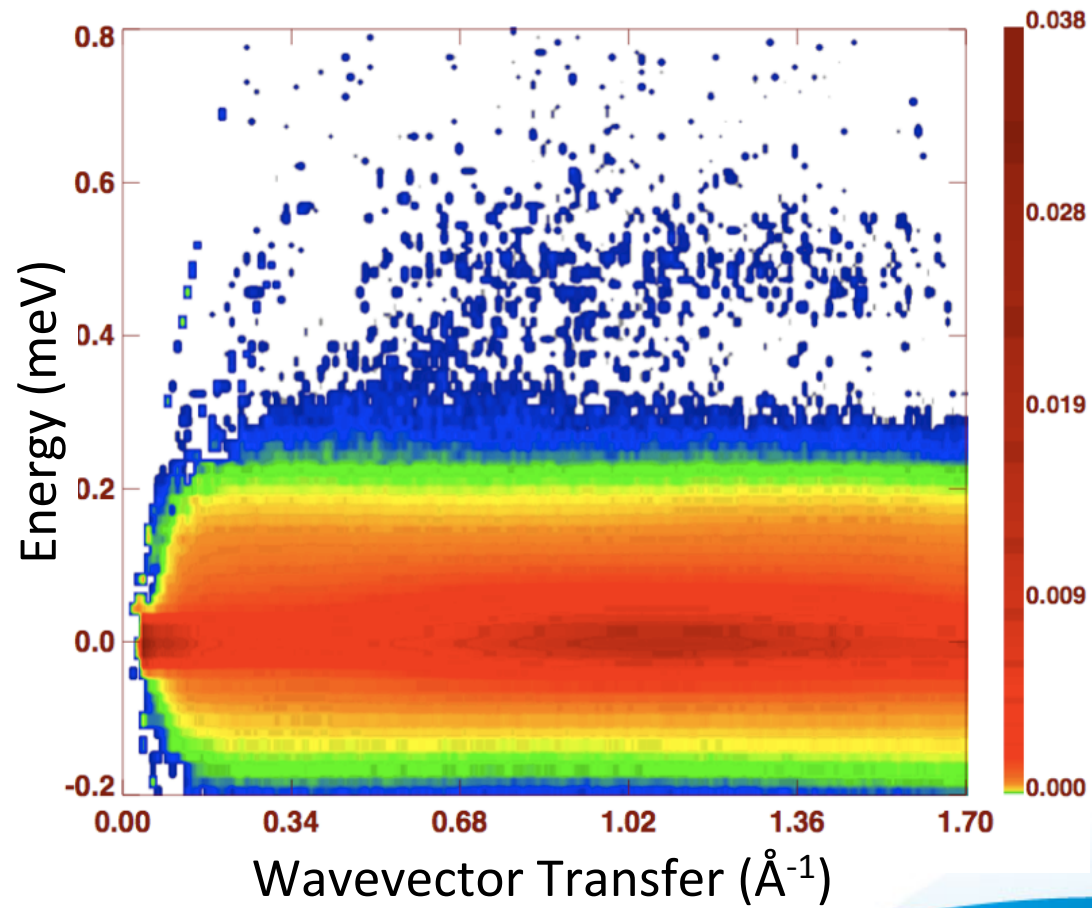




# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

1.2 K

FT of spin spin autocorrelation.



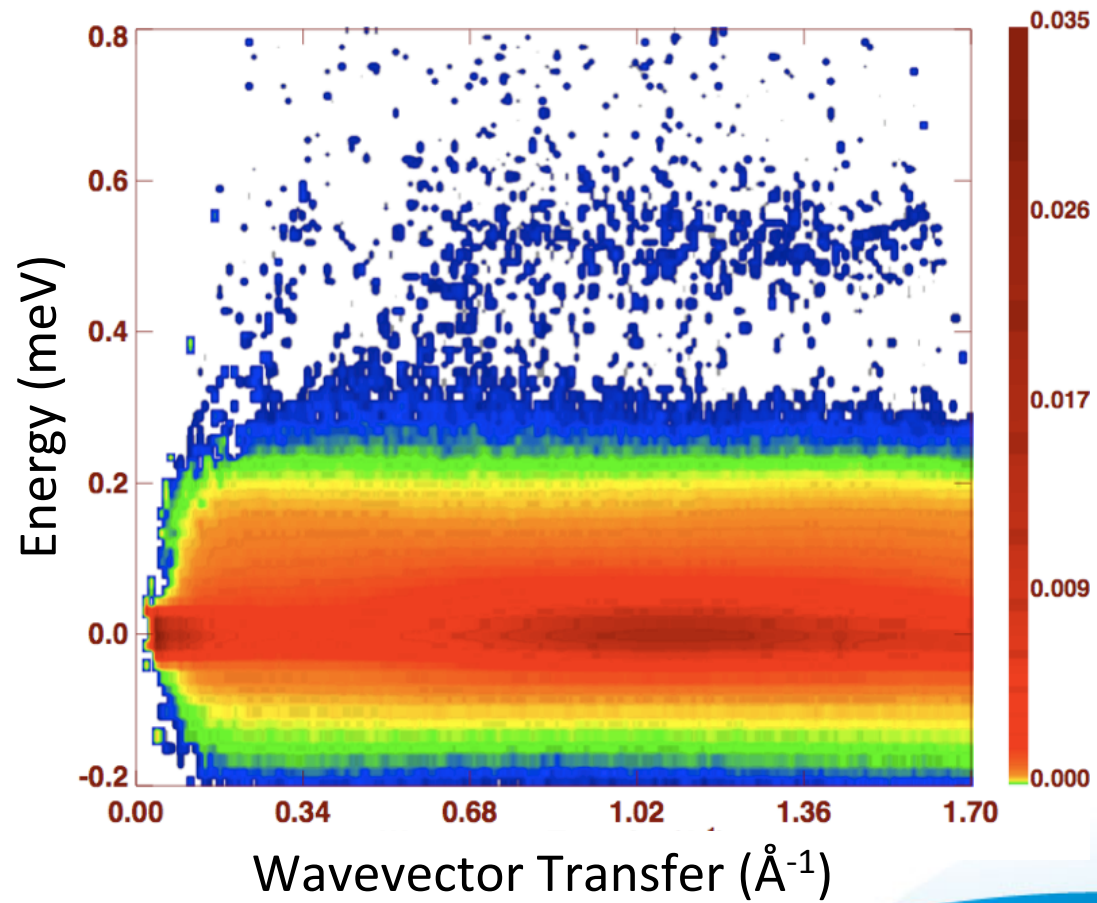




# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

0.9 K

FT of spin spin autocorrelation.

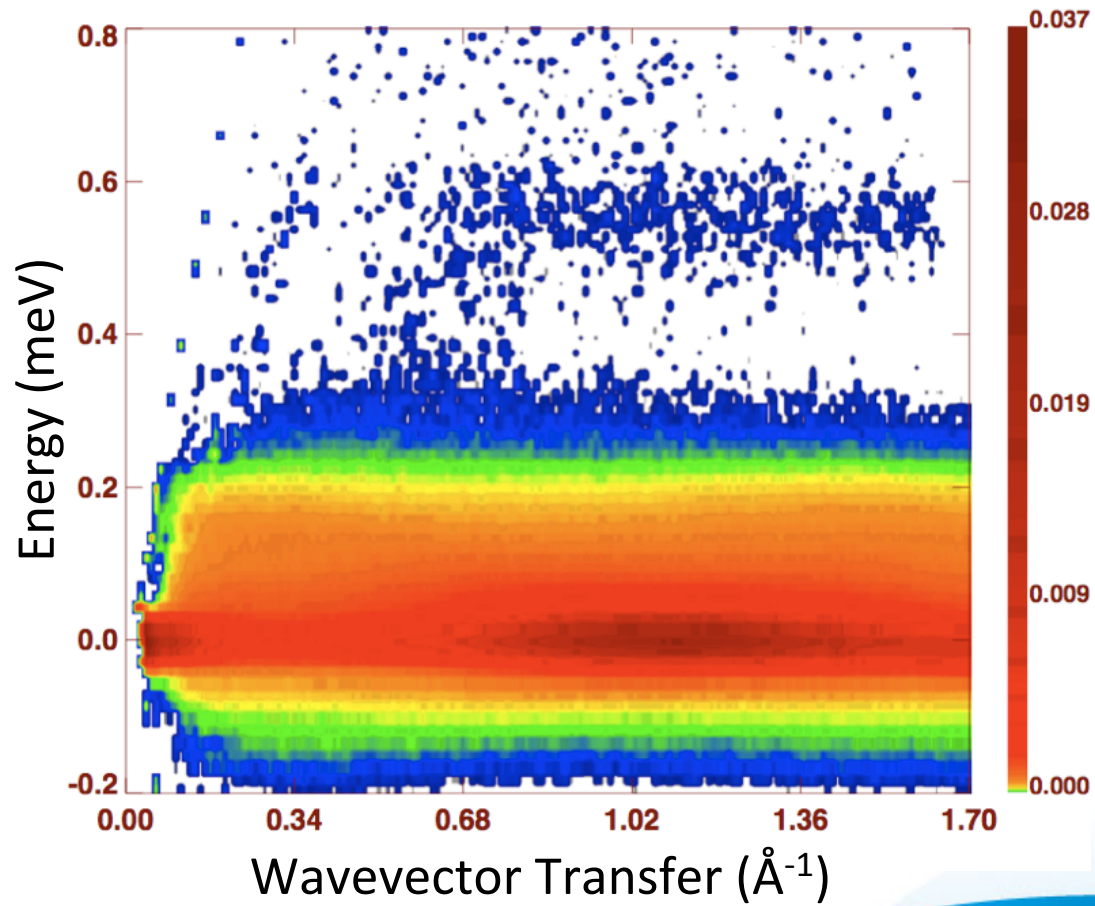




# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

0.6 K

FT of spin spin autocorrelation.

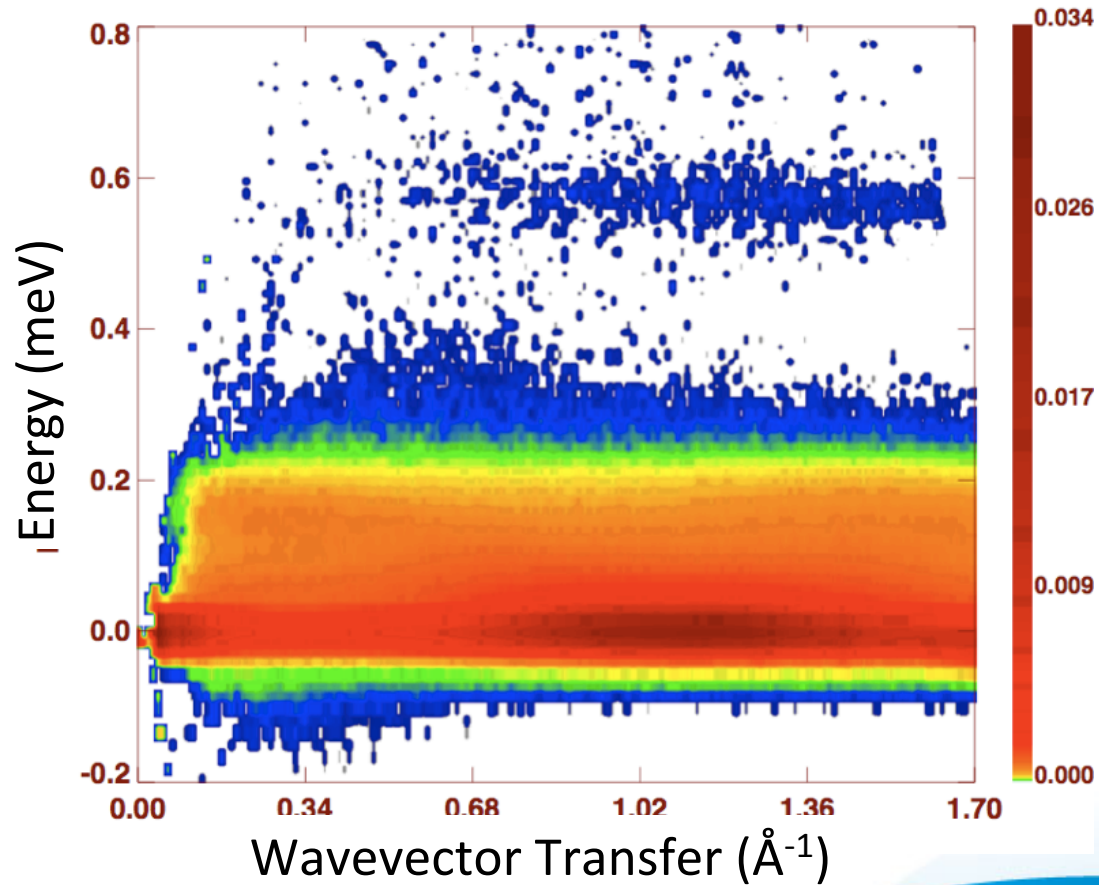




# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

0.25 K

FT of spin spin autocorrelation.

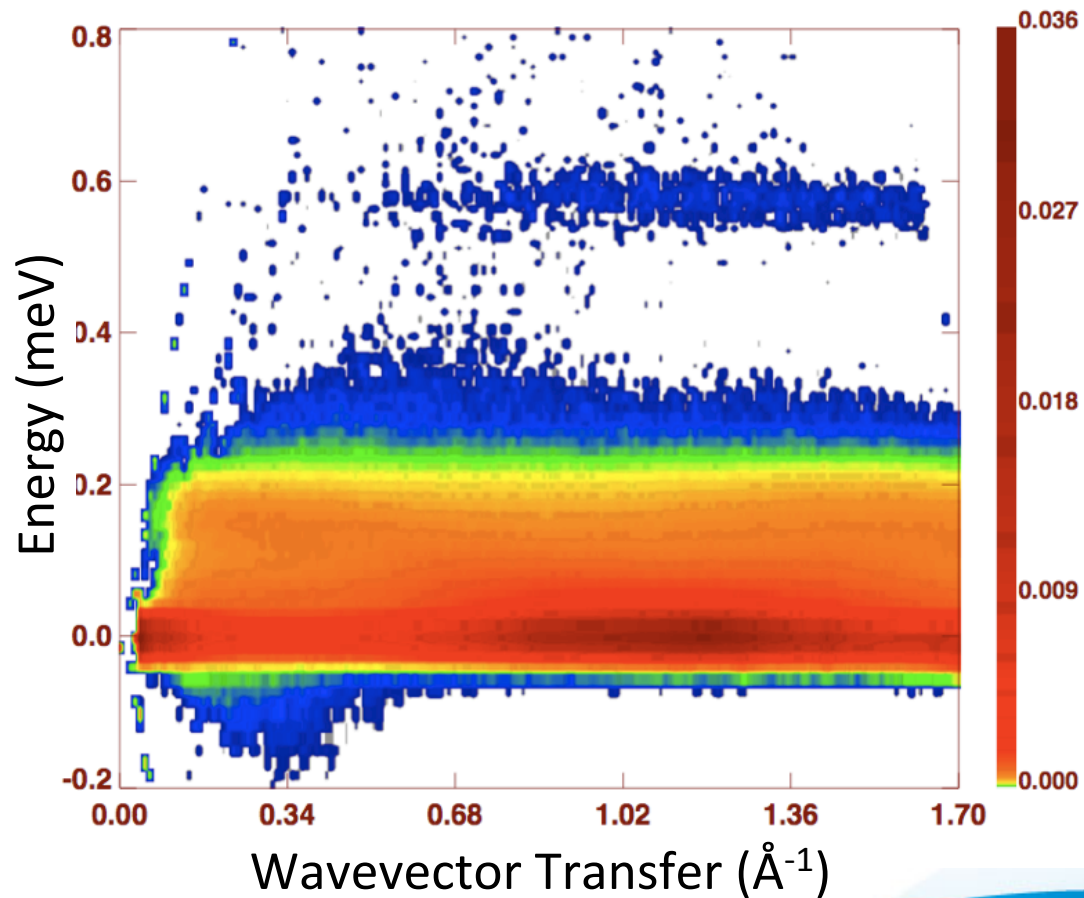




# Inelastic neutron scattering (IN5) ( $\Delta E = 50 \mu\text{eV}$ )

0.06 K (Below 0.14 K transition)

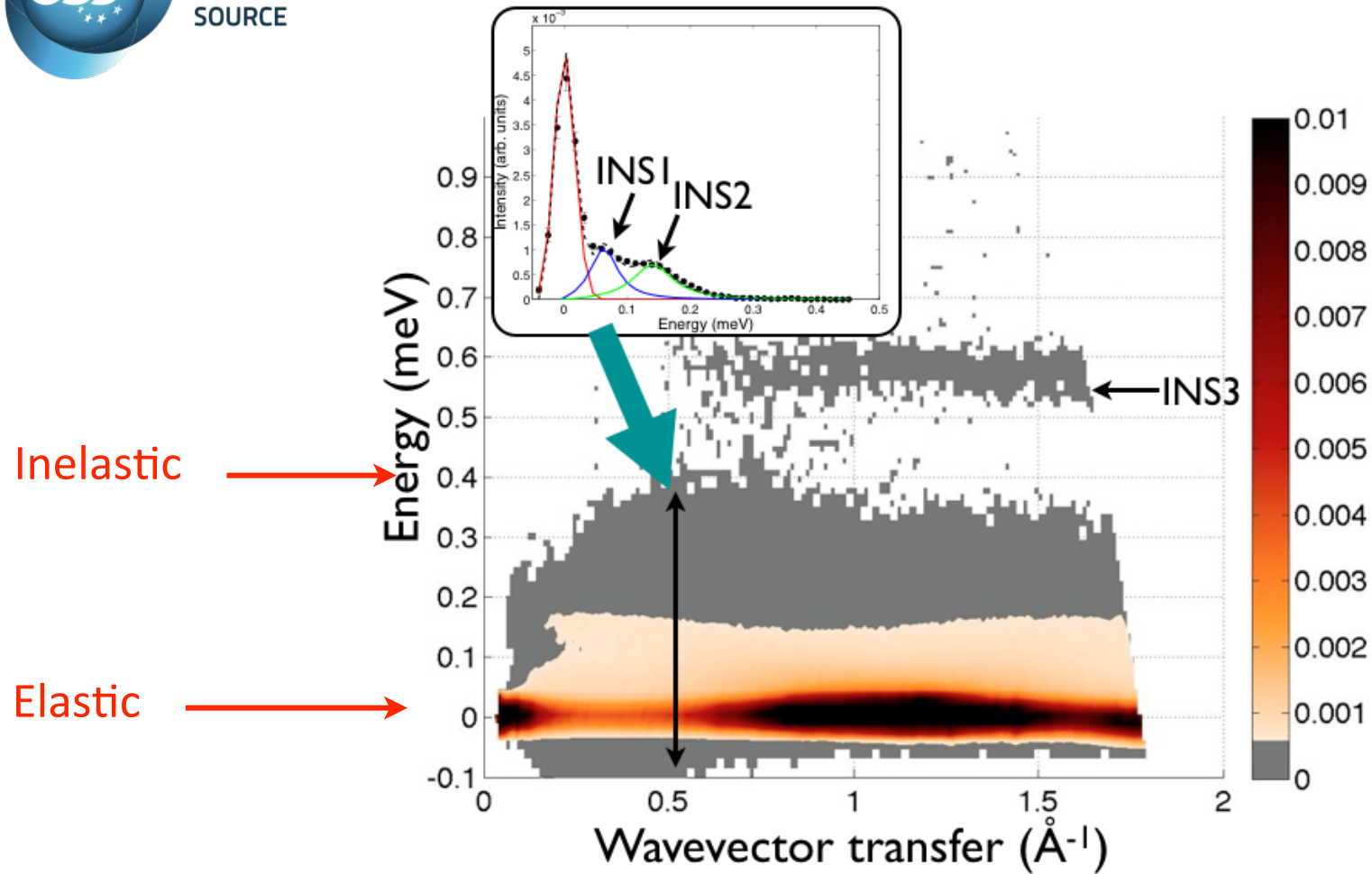
FT of spin spin autocorrelation.



1st inelastic neutron scattering data on a frustrated 3D Kagome compound.



After some analysis

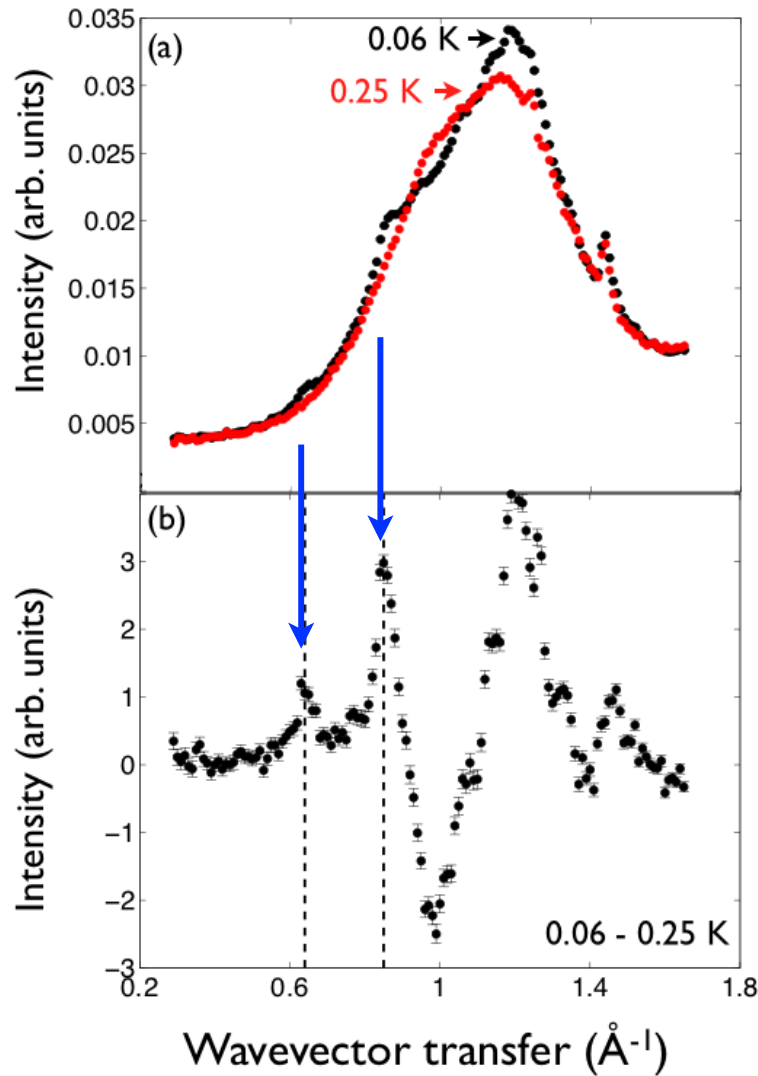


3 - gapped non-dispersive excitations  
INS1 = 0.04, INS2 = 0.14 and INS3 = 0.58 meV





# Elastic scattering



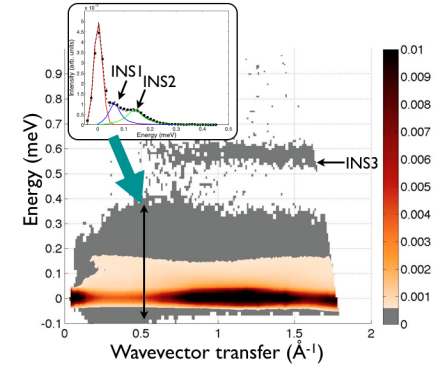
82 % of the total scattering intensity

Incommensurate order  
( $T < 0.14 \text{ K}$ ) within elastic line  
Frozen order is indeed frozen (within resolution)  
No spin wave excitations.



# Inelastic scattering

Non-dispersive:  
crystal fields?  
Local vibrational excitations?



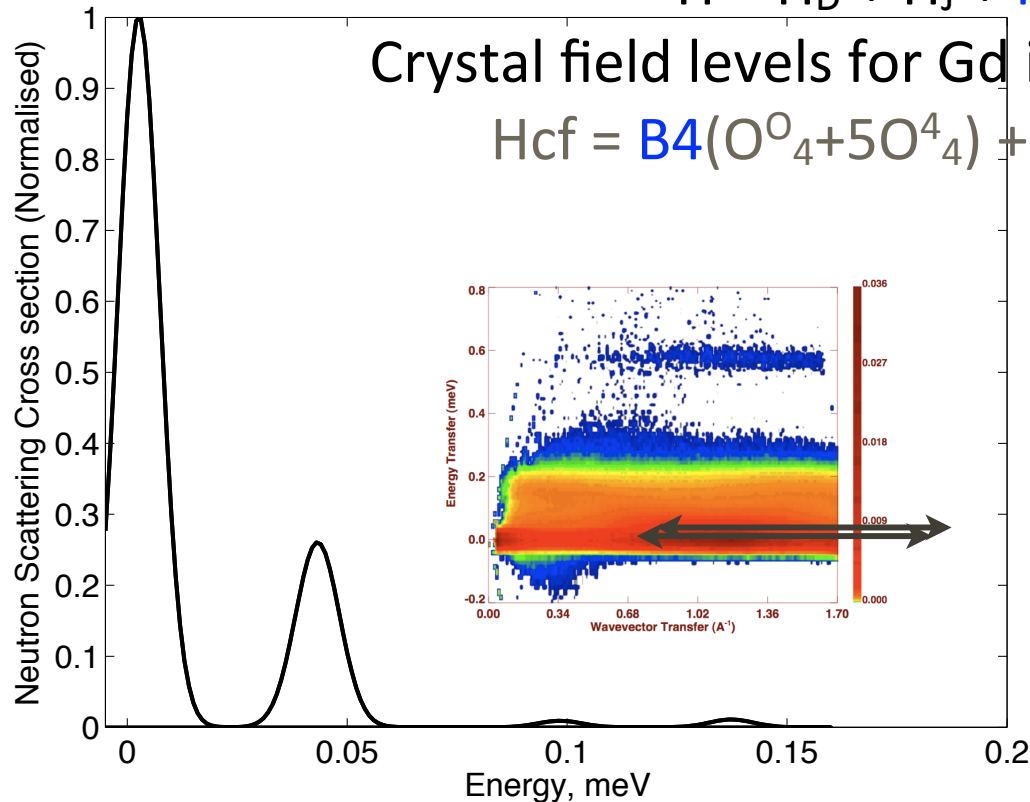
Newman et al. J. Phys. C. (8),1862, 1975

$$H = H_D + H_J + H_{CF} + H_Z$$

Crystal field levels for Gd in garnet structure.

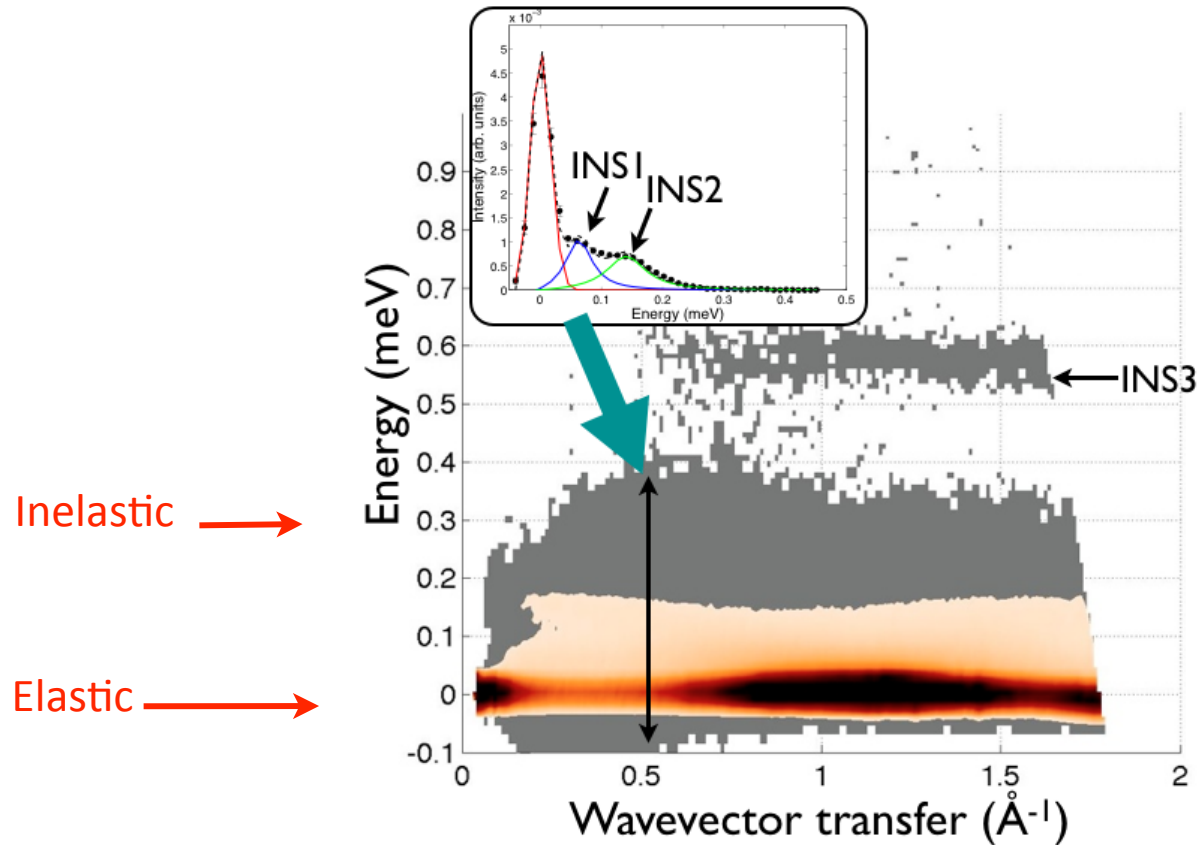
$$H_{cf} = B_4(O_4^0 + 5O_4^4) + B_6(O_6^0 - 21O_6^6)$$

Cannot be reconciled with bulk measurements





## Data Analysis



Q-dependence - integrate over inelastic line at each position of Q.

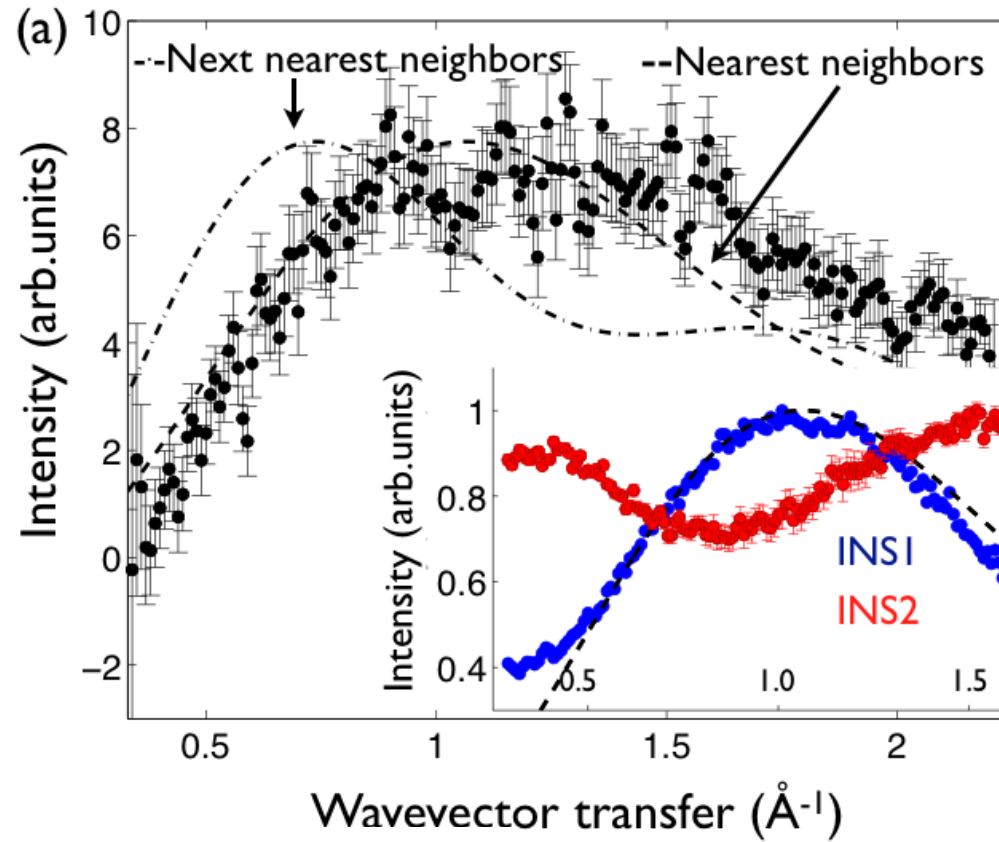
- Q - dependence 60 mK.
- Temperature dependence.





# Inelastic scattering 60 mK - possible origin

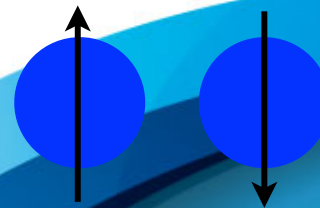
## AF short range correlations



INS1 ✓  
INS2 ✗  
INS3 ✓

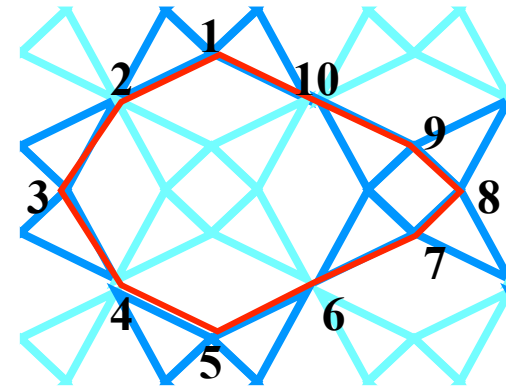
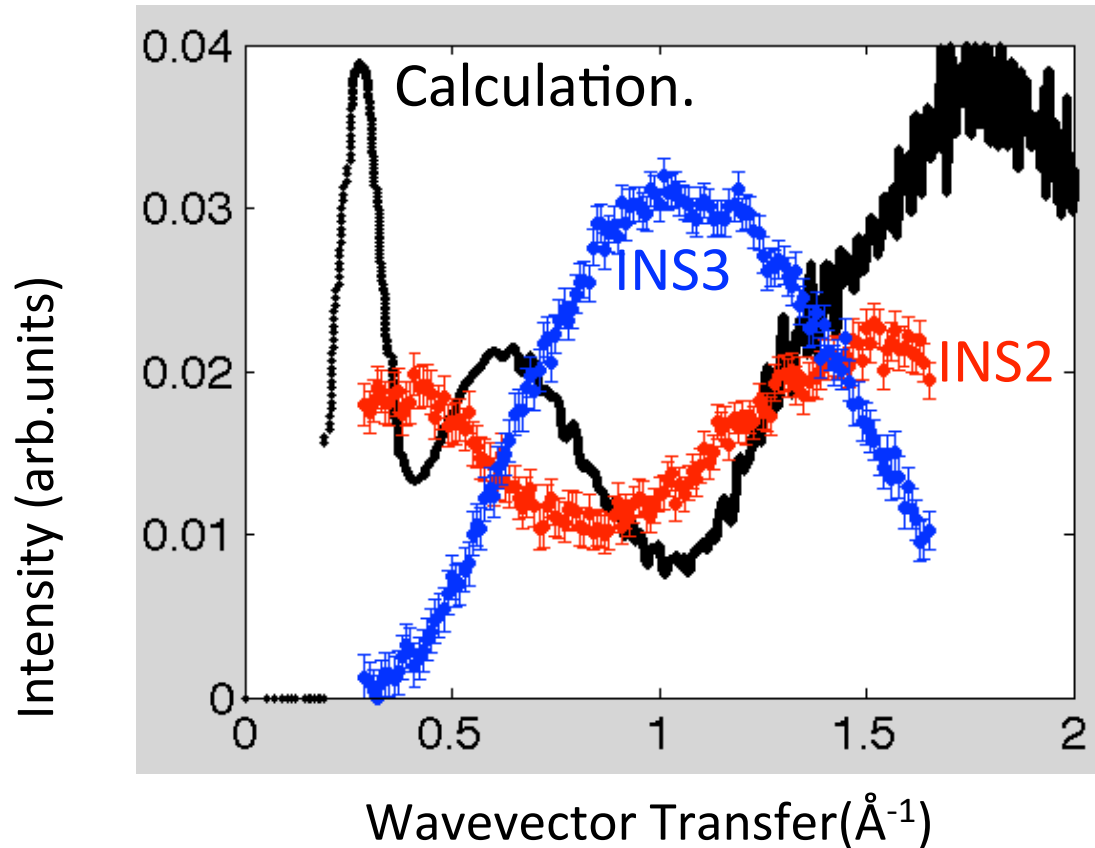
$$\frac{d^2\sigma}{d\Omega d\omega} \propto A(T) F^2(Q) \left[ 1 - \frac{\sin(Qd)}{Qd} \right]$$

$J_{NN} \sim 1.68 \text{ K}$





## Inelastic scattering 60 mK - 10 spin cluster?



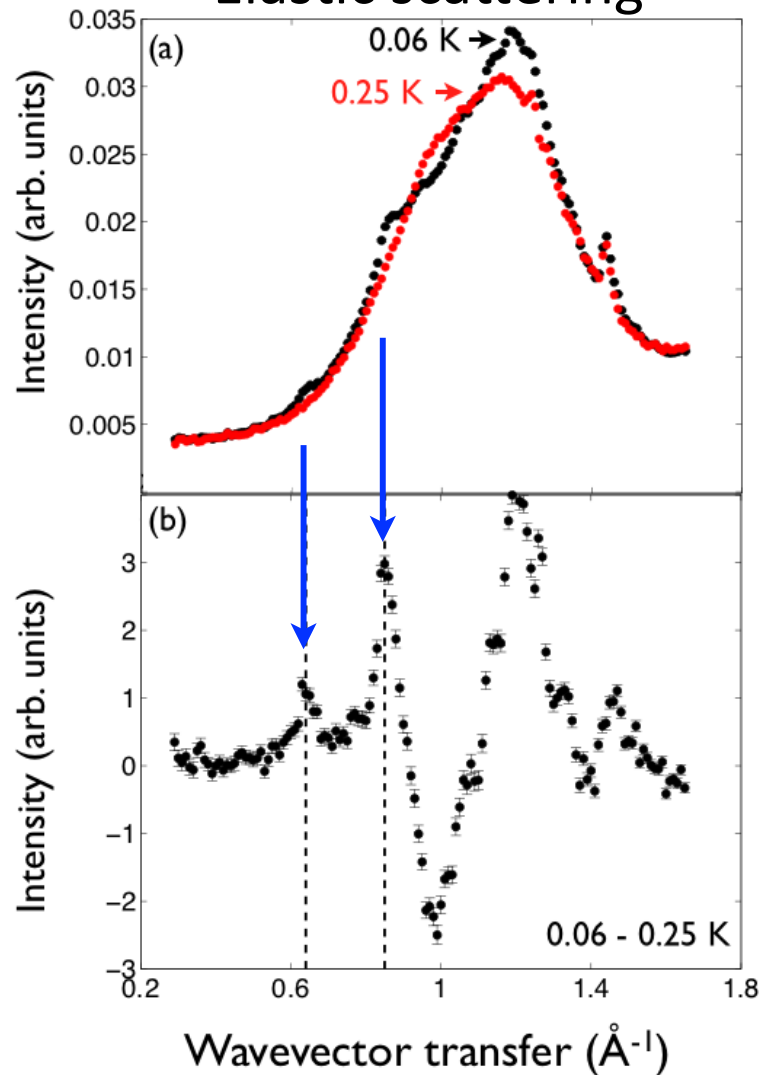
INS3: Not a 10 spin cluster  
INS2: close to 10 spin cluster



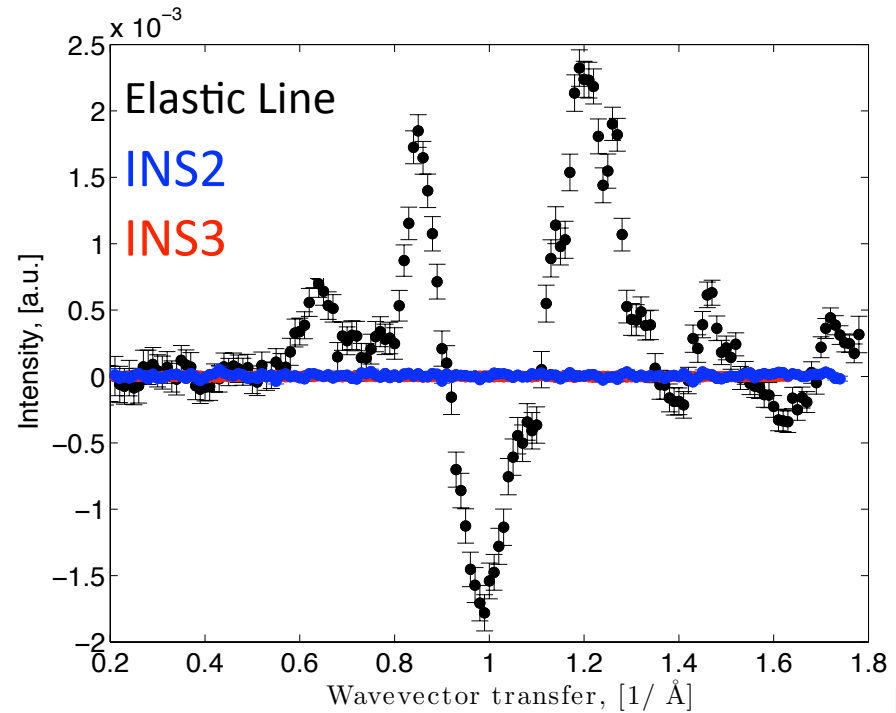
# Inelastic scattering : Temperature dependence

60 mK - 250 mK (Across 0.14 K transition)

### Elastic scattering



### Inelastic scattering

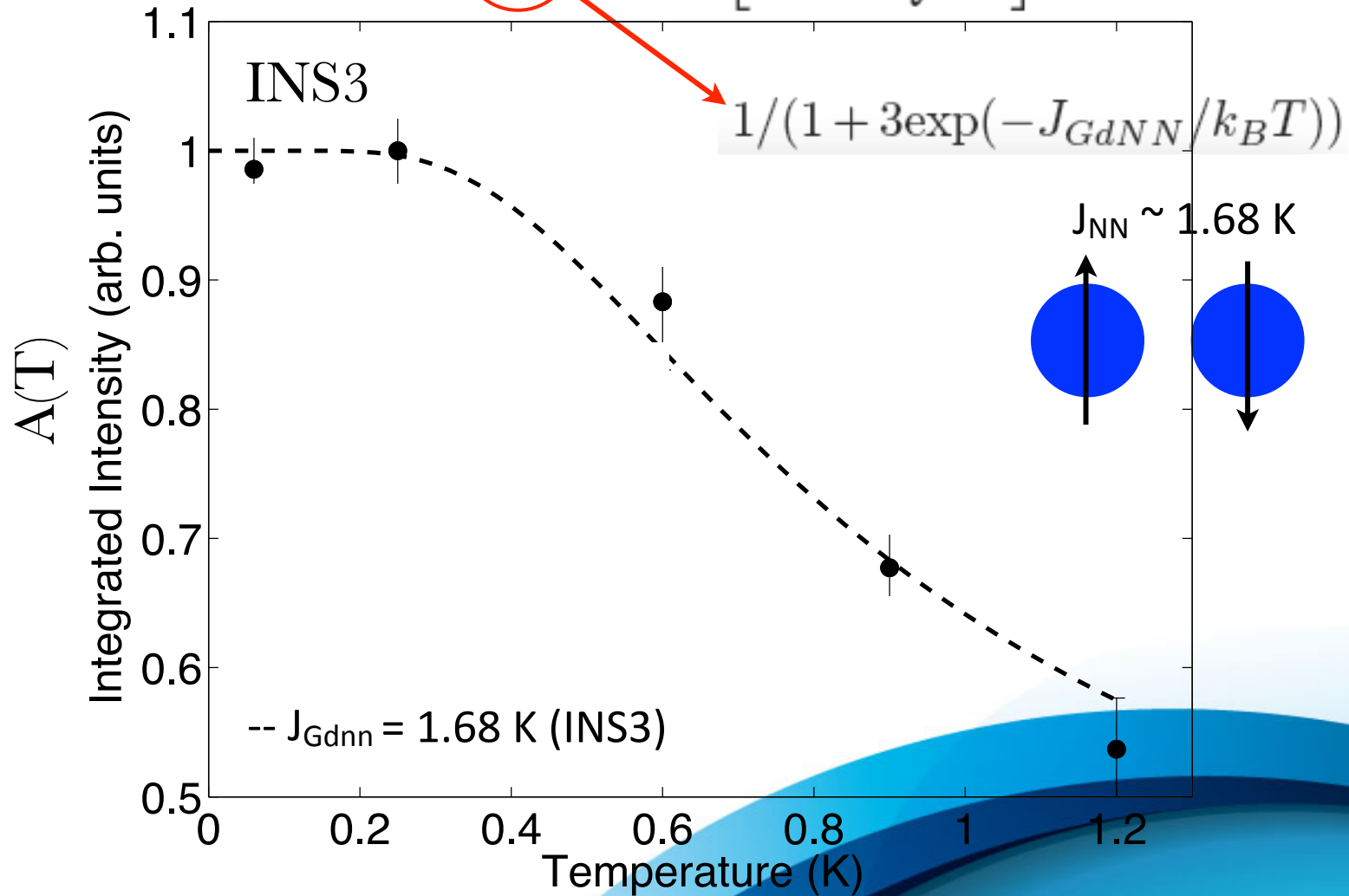


Excitations not affected by longer range order



# Inelastic scattering :Temperature dependence (Thermal distribution of singlet to triplet excitations)

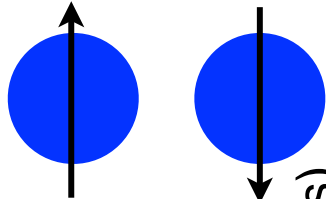
$$\frac{d^2\sigma}{d\Omega d\omega} \propto A(T) F^2(Q) \left[ 1 - \frac{\sin(Qd)}{Qd} \right]$$



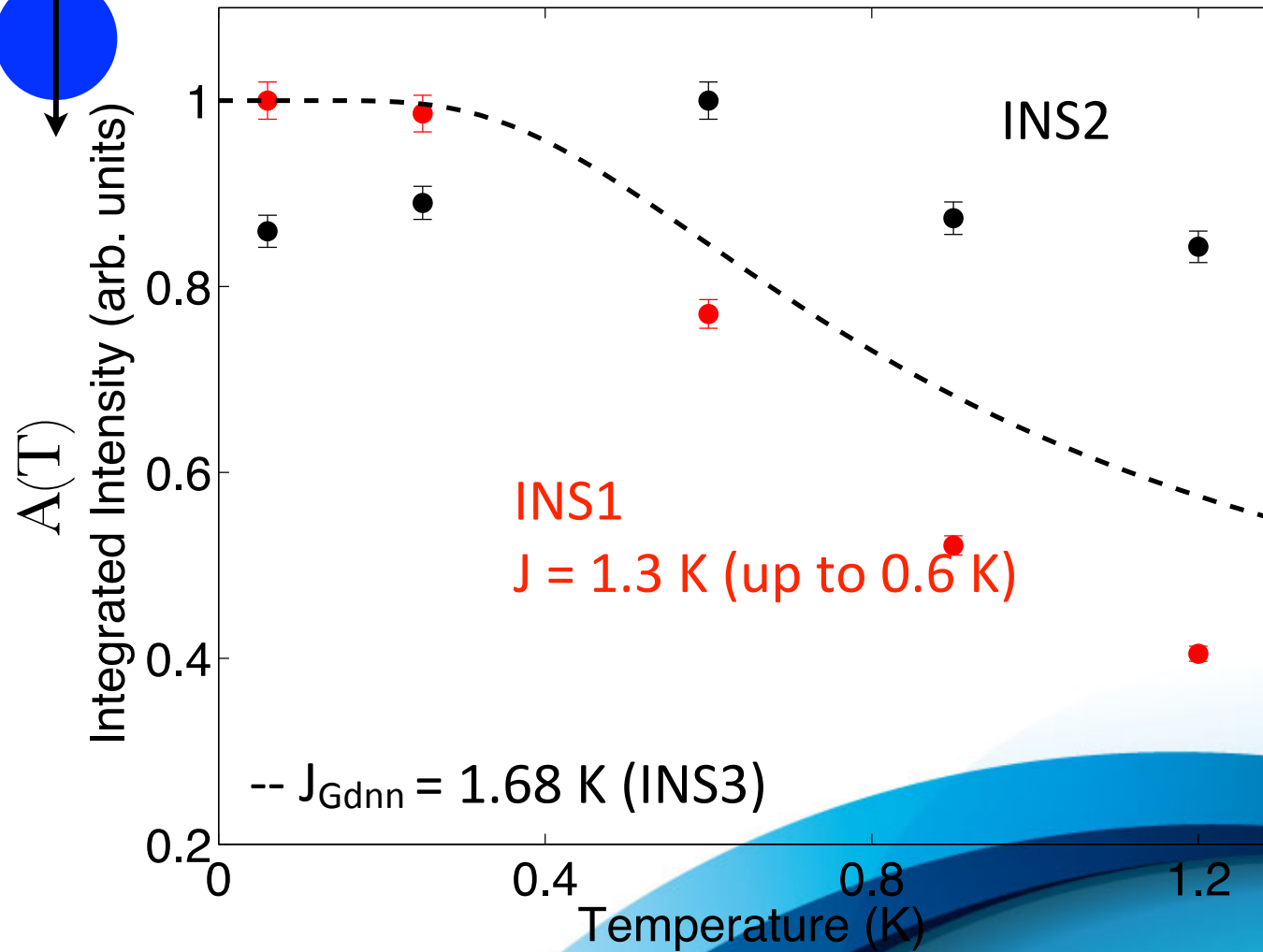
# Inelastic scattering : Temperature dependence

(Thermal distribution of singlet to triplet excitations)

$J_{NN} \sim 1.68$  K



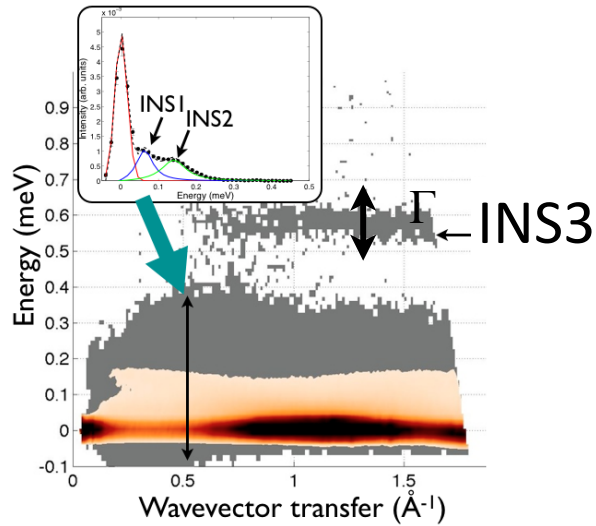
$$\frac{d^2\sigma}{d\Omega d\omega} \propto A(T)F^2(Q) \left[ 1 - \frac{\sin(Qd)}{Qd} \right]$$



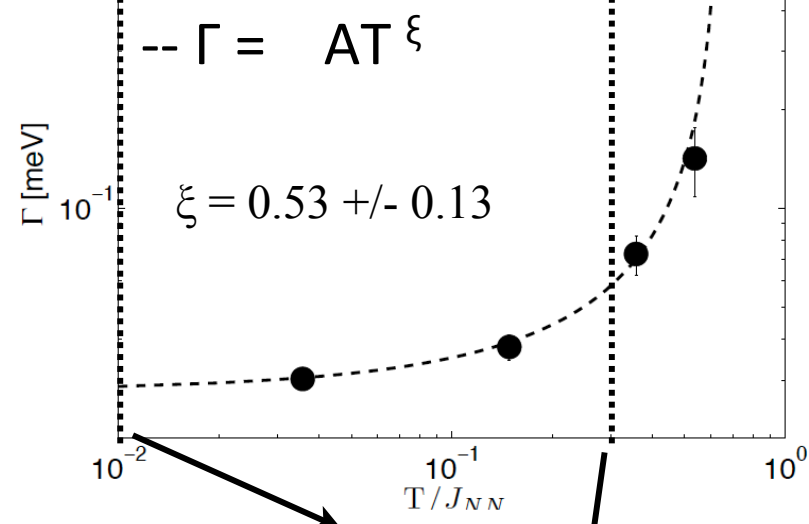


# Inelastic scattering : Temperature dependence

## INS3 Inelastic Width

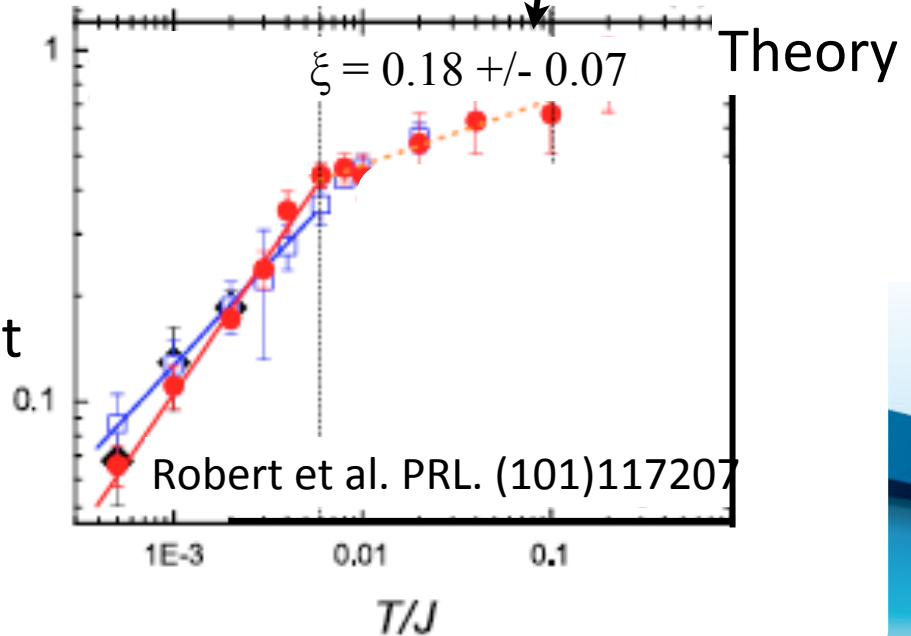


Dynamical spin-pair correlations Experimental



Scattering follows algebraic law  
 $\Gamma = AT^\xi$  but exponent not equivalent

Robert only calculated  
 J Near neighbour

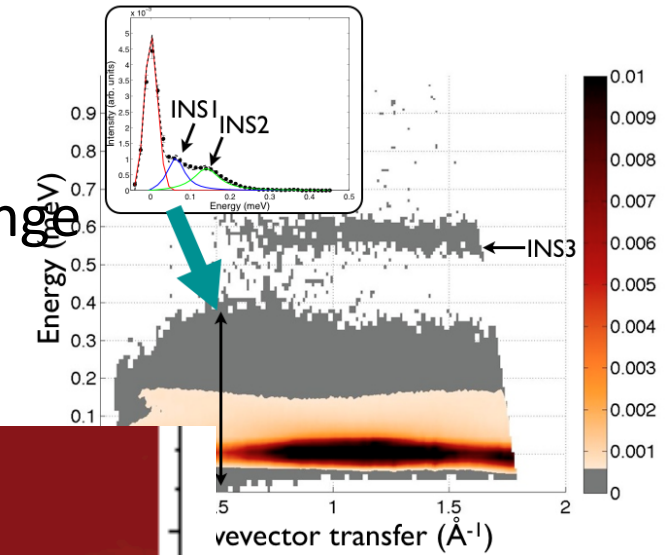
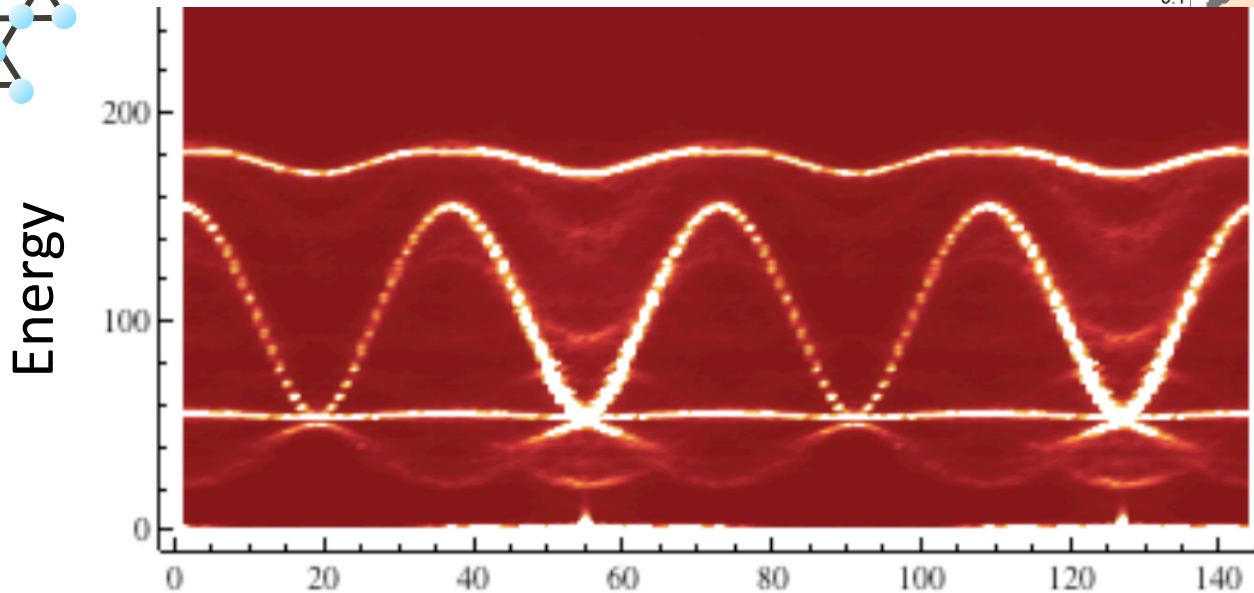
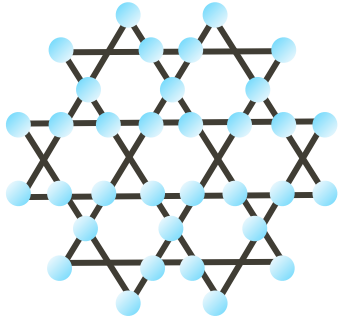




# Further Theory for Kagome

## B. Canals

Near neighbour and dipole exchange interactions only  
Are these soft modes?



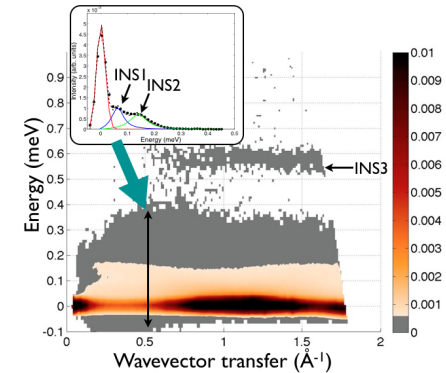
- Always 3 clear modes
- Lowest always flat
- Second always dispersive
- Appears that  $J_1$  only drives model into a coplanar gapped modes (+D = gapped)
- Role of temperature is not important.

Spatial Correlation



## Some Conclusions

- No continuum of excitations in 3D Kagome.
  - Soft modes exist in a zero field disordered state
  - Data consistent with AF short range exchange (3 spin)
  - Excitations affected by  $J_{nn}$  - SRO.
  - Long range order does not perturb excitations.
- 
- Are these soft modes lifted by anisotropy?
  - Must include Anisotropic dipole interactions.
  - No spin waves?



PHYSICAL REVIEW B **82**, 174408 (2010)



### **Spin dynamics in the hyperkagome compound $Gd_3Ga_5O_{12}$**

P. P. Deen,<sup>1</sup> O. A. Petrenko,<sup>2</sup> G. Balakrishnan,<sup>2</sup> B. D. Rainford,<sup>3</sup> C. Ritter,<sup>1</sup> L. Capogna,<sup>4</sup> H. Mutka,<sup>1</sup> and T. Fennell<sup>1</sup>



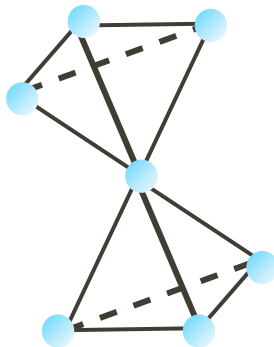


# Frustration $\Rightarrow$ zero point entropy $\Rightarrow$ degeneracy

Universal dynamic behaviour?

It would appear not.

Pyrochlore

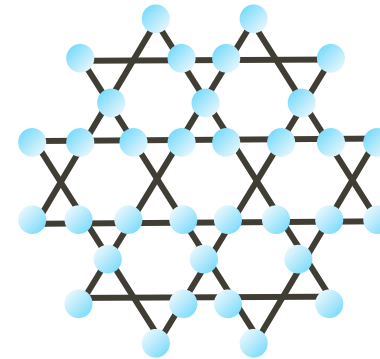


Dynamic gapless excitations to 0 K.

$$\Delta E \sim T.$$

Depends on T not J.

Kagome/3D Kagome



3 Gapped excitations

Algebraic relation for  $\Delta E$ .

Weak temperature dependence

Dynamic range - meV - picoeV

Thank you.



EUROPEAN  
SPALLATION  
SOURCE

## Spectroscopy into the future

The study of dynamics in condensed matter with inelastic neutron scattering provides one of the most exacting tests of the understanding of the microscopic origin of the material properties, particularly when combined with powerful computer modelling techniques now being pioneered.

[LET \(ISIS\) Website](#)



## Scientific & Sociatal Impact

### Magnetism

Electron -phonon coupling

Quantum phase transition

New classes of quantum criticality

Frustrated magnetism

Electron, magnetic and SC order

Strongly electron correlations

Giant/Collosal Magnetoresistance

Molecular magnetism

### Chemical Physics

Confinement

Collective behaviour in liquids and glasses

Hydrogen storage

Molecular dynamics of living cells

Targeted drug delivery



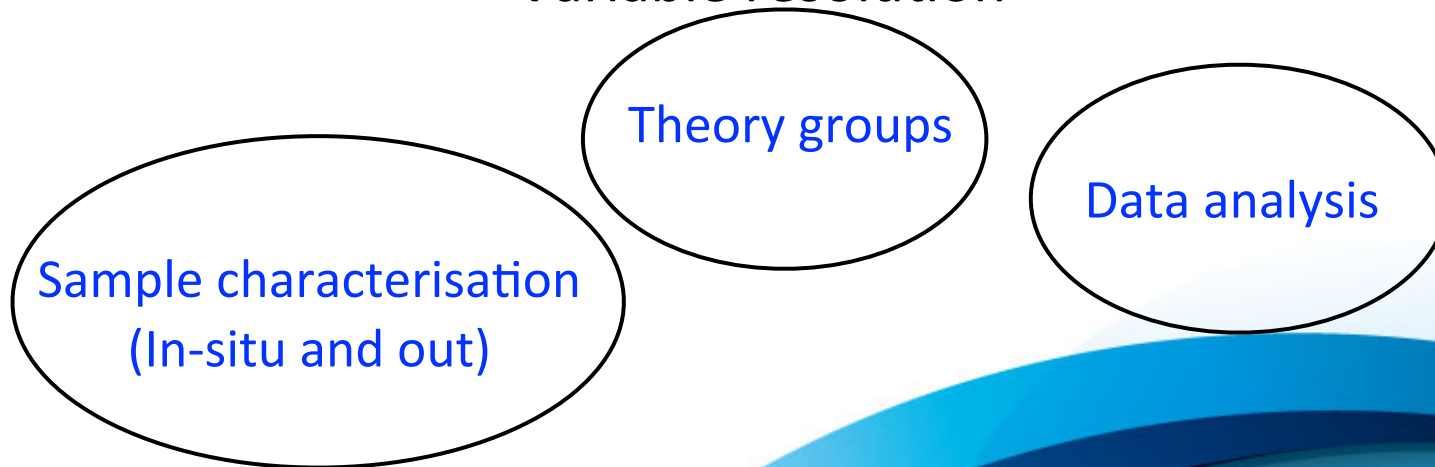
## Spectroscopy into the future

Weak continuum scattering

Polarisation analysis

A few examples

Variable resolution





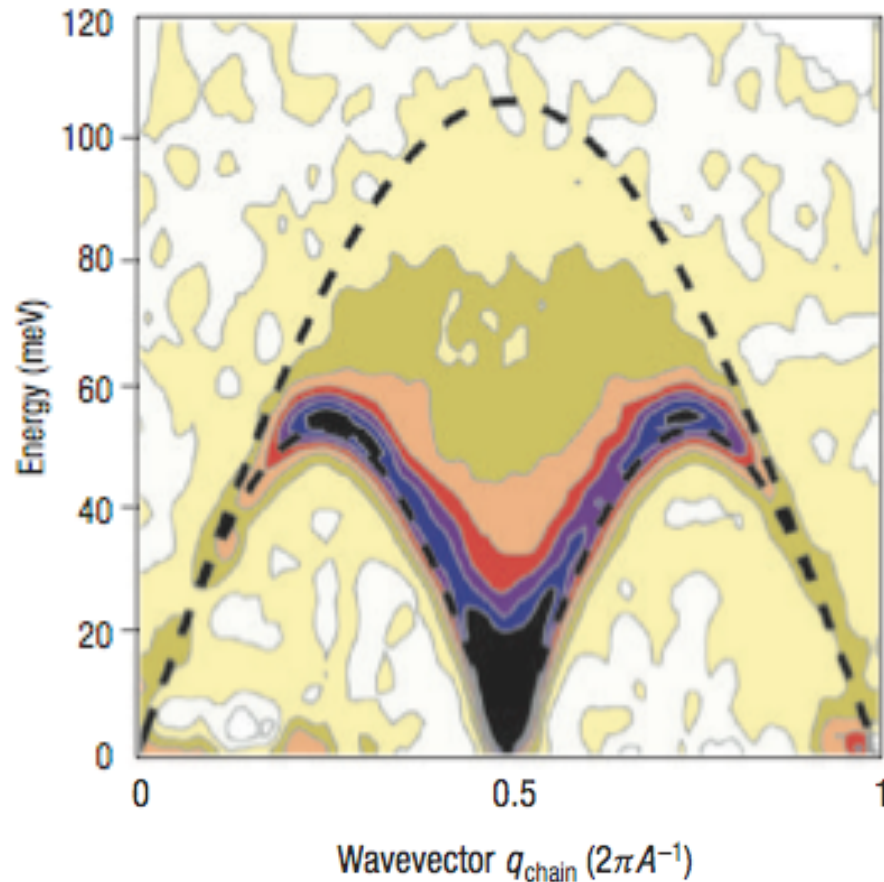
EUROPEAN  
SPALLATION  
SOURCE

## (I) Large areas of $S(Q, \omega)$

Coming from the magnetism angle

Total  $S(Q, \omega)$   $\text{KCuF}_3$

B.Lake et al. Nature Materials (2005)



Weakly coupled 1D spin chains  
-- predicted multi-spinon continuum.

model for carbon nanotube

Dynamic behaviour close to quantum  
criticality

Scattering extends over a broad  
region of  $S(Q, \omega)$ .

Variable resolution.

Valuable and complex physics sits within the  
“background”.



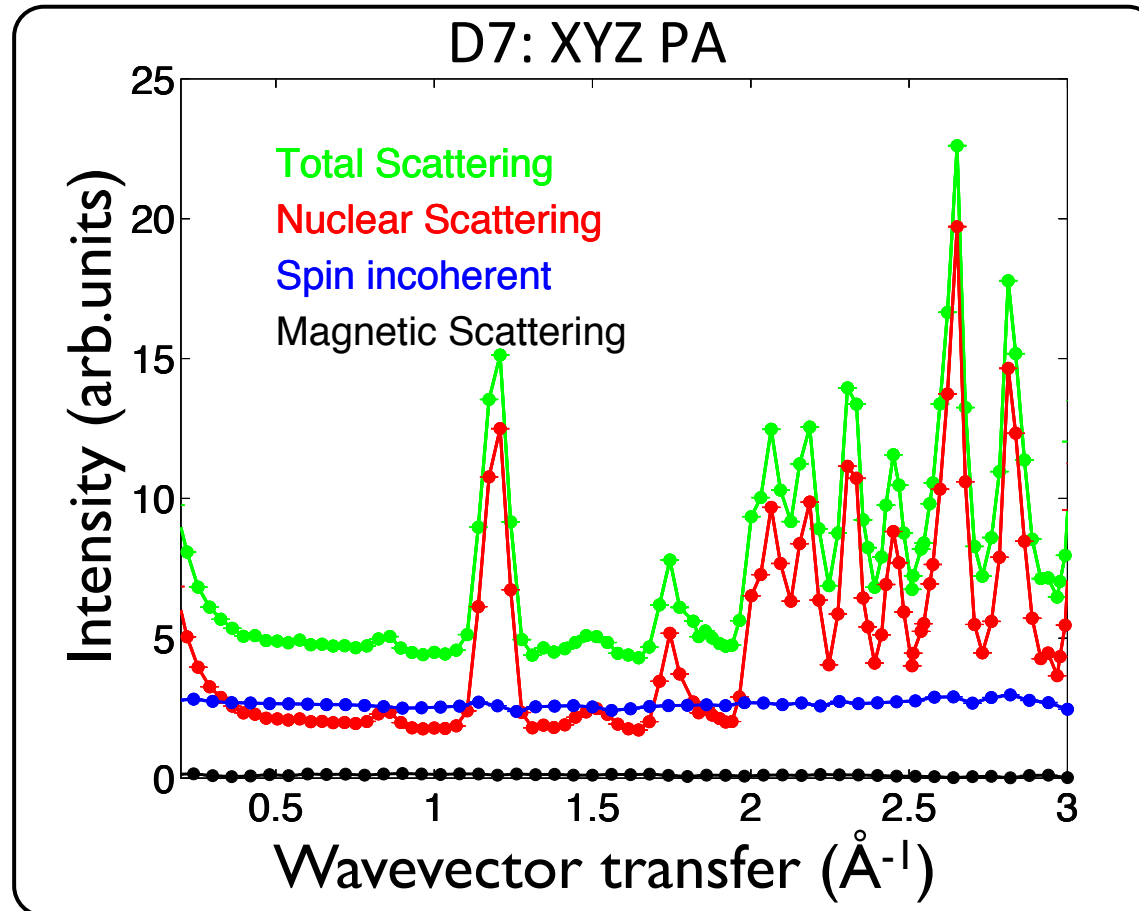
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SPALLATION  
SOURCE

## (II) PA

Pair Correlations, Short Range Order and Dispersive Excitations in the  
Quasi-Kagome Quantum Magnet Volborthite

G. J. Nilsen,<sup>1,2,3,\*</sup> F. C. Coomer,<sup>3,†</sup> M. A. de Vries,<sup>4</sup> J. R. Stewart,<sup>5</sup> P. P. Deen,<sup>6</sup> A. Harrison,<sup>3,7</sup> and H. M. Rønnow<sup>1</sup>

[arXiv:1001.2462](https://arxiv.org/abs/1001.2462)



$S = 1/2$  Kagome Heisenberg antiferromagnet  
(True quantum spin liquid?)

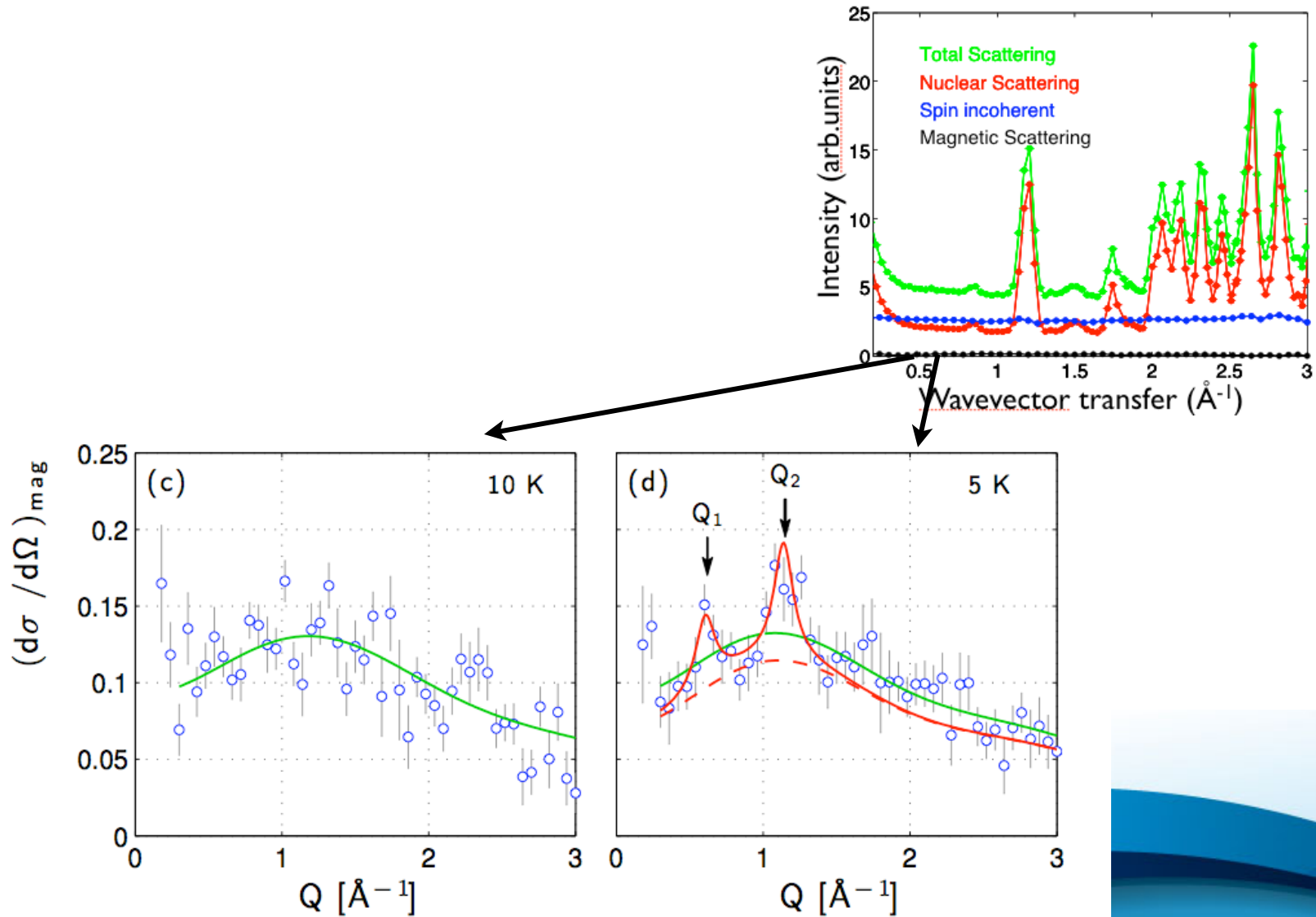




## (II) PA

### Pair Correlations, Short Range Order and Dispersive Excitations in the Quasi-Kagome Quantum Magnet Volborthite

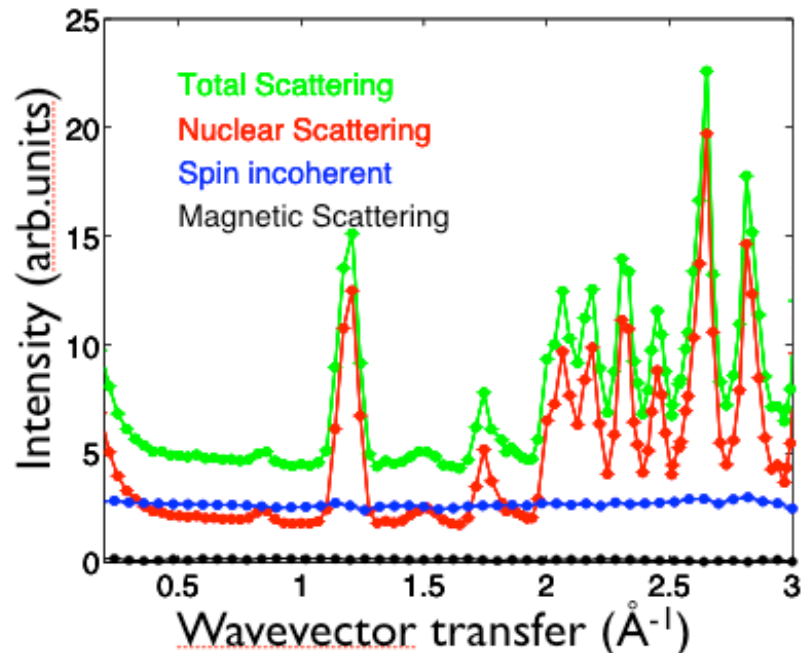
G. J. Nilsen,<sup>1,2,3,\*</sup> F. C. Coomer,<sup>3,†</sup> M. A. de Vries,<sup>4</sup> J. R. Stewart,<sup>5</sup> P. P. Deen,<sup>6</sup> A. Harrison,<sup>3,7</sup> and H. M. Rønnow<sup>1</sup>





## Polarisation analysis

Polarisation analysis is imperative in the study of magnetic order.



The unambiguous observation of magnetic continuum scattering is highly nontrivial, because of the difficulty of isolating weak magnetic continua from background and phonon contributions. **Definite proof that additional scattering is magnetic can only come from polarisation analysis.**

N. B. Christensen et al. PNAS, vol. 104, p. 15264, 2007.

However other communities are less enthusiastic!

Loss of flux is not appreciated



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SOURCE

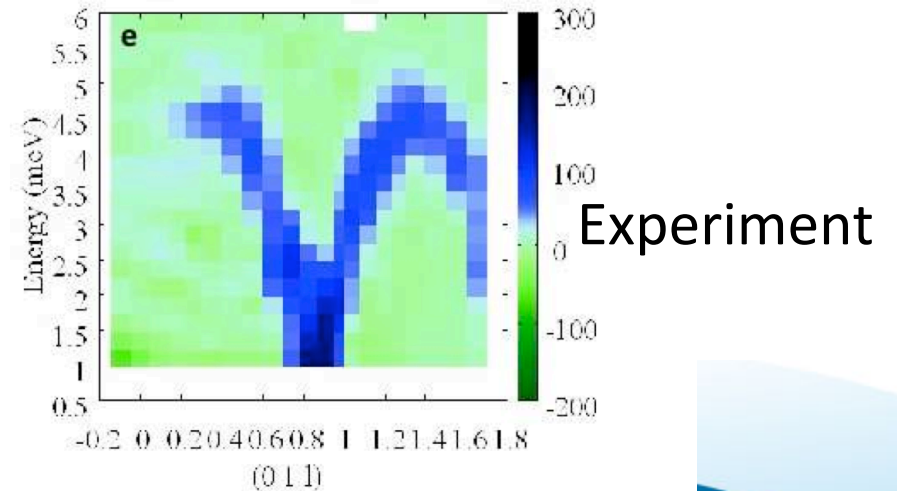
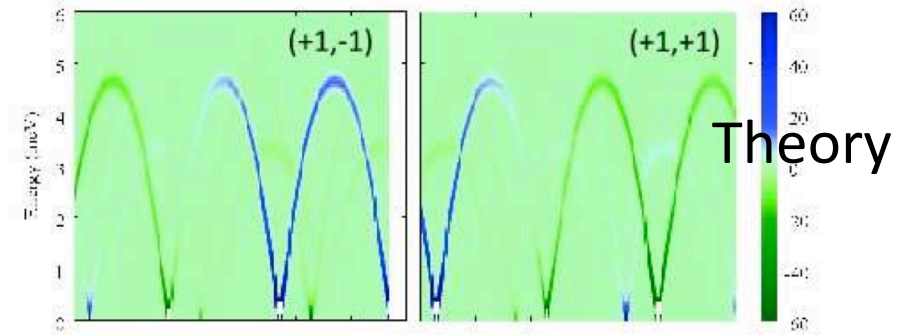
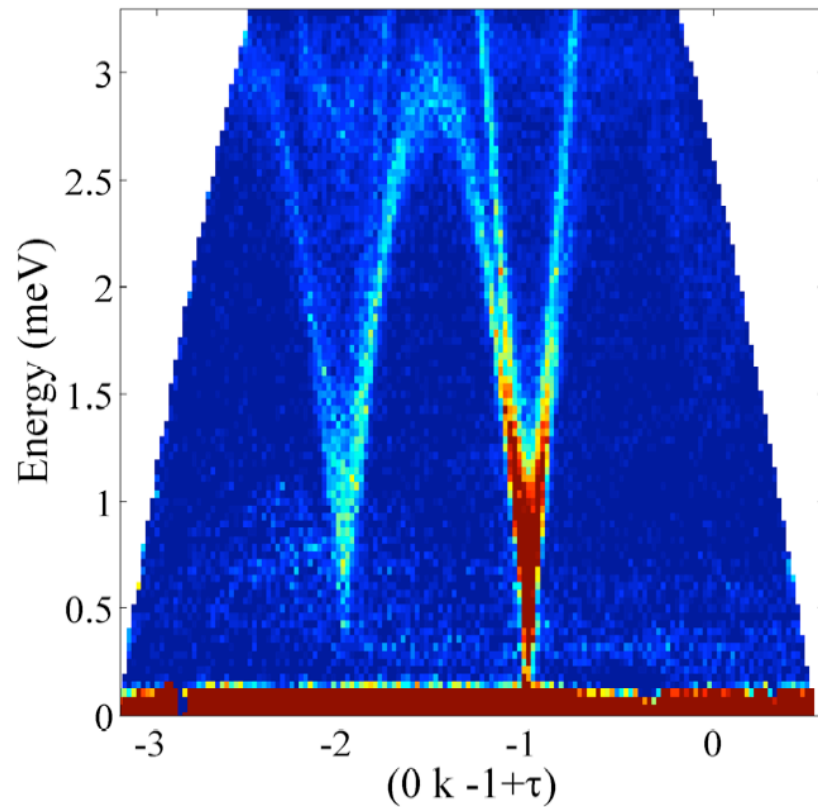
## (II)(b) PA

When magnetic chirality gets excited: spin waves in  $\text{Ba}_3\text{NbFe}_3\text{Si}_2\text{O}_{14}$

M. Loire,<sup>1</sup> V. Simonet,<sup>1</sup> S. Petit,<sup>2</sup> K. Marty,<sup>1,3</sup> P. Bordet,<sup>1</sup> P. Lejay,<sup>1</sup> J. Ollivier,<sup>4</sup> M. Enderle,<sup>4</sup> P. Steffens,<sup>4</sup> E. Ressouche,<sup>5</sup> A. Zorko,<sup>6</sup> and R. Ballou<sup>1,\*</sup>

### IN20 (CRYOPAD) - PA

### IN5 (ILL)

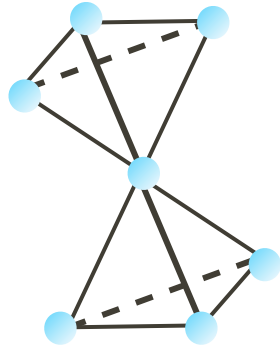


<http://horace.isis.rl.ac.uk/>

**Magnetic Coulomb Phase in the Spin Ice  $\text{Ho}_2\text{Ti}_2\text{O}_7$**

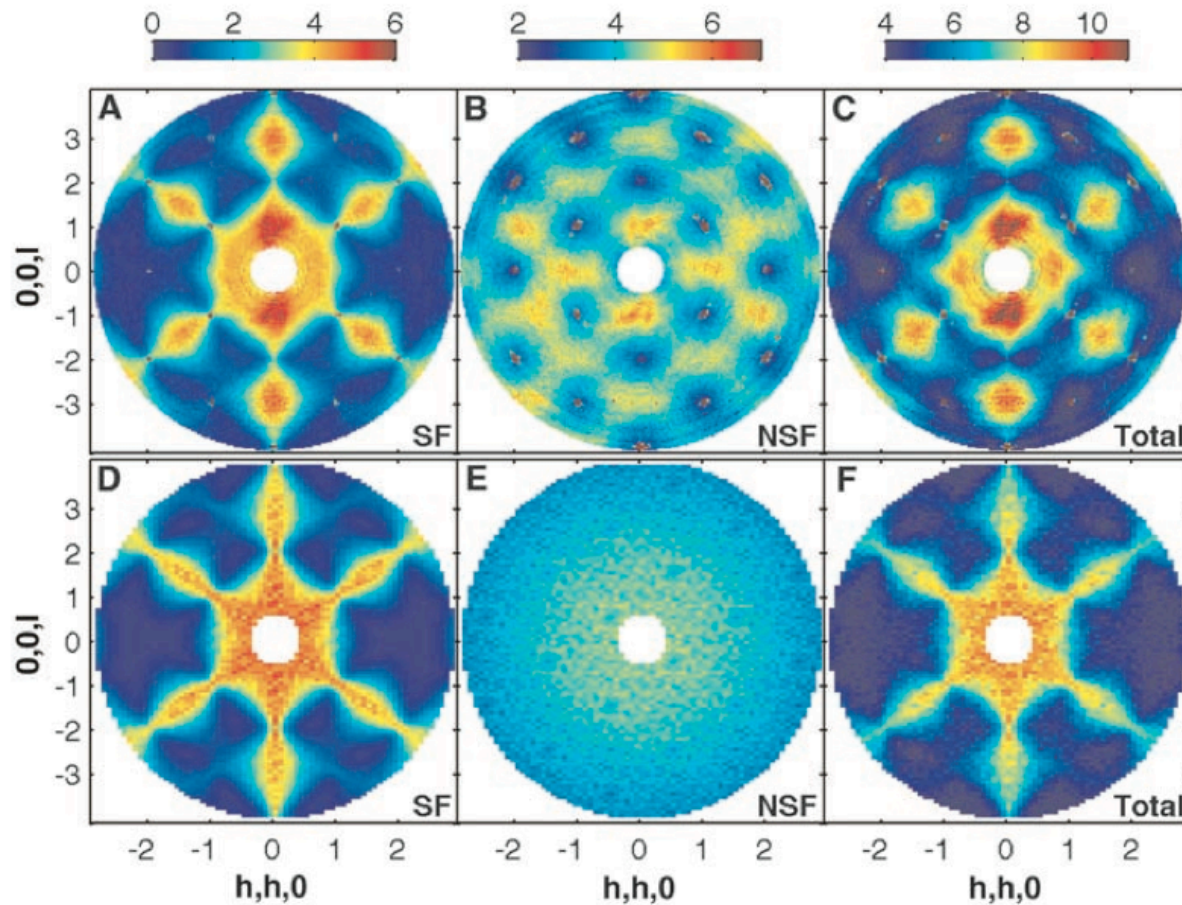
T. Fennell<sup>1,\*</sup>, P. P. Deen<sup>1</sup>, A. R. Wildes<sup>1</sup>, K. Schmalzl<sup>2</sup>, D. Prabhakaran<sup>3</sup>, A. T. Boothroyd<sup>3</sup>, R. J. Aldus<sup>4</sup>,  
D. F. McMorrow<sup>4</sup> and S. T. Bramwell<sup>4</sup>

Published Online 3 September 2009  
*Science* 16 October 2009:  
Vol. 326 no. 5951 pp. 415–417



Experiment

Theory

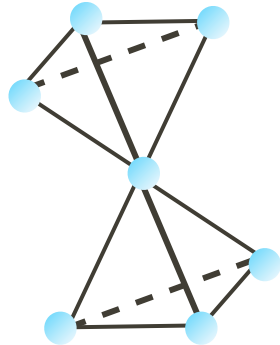




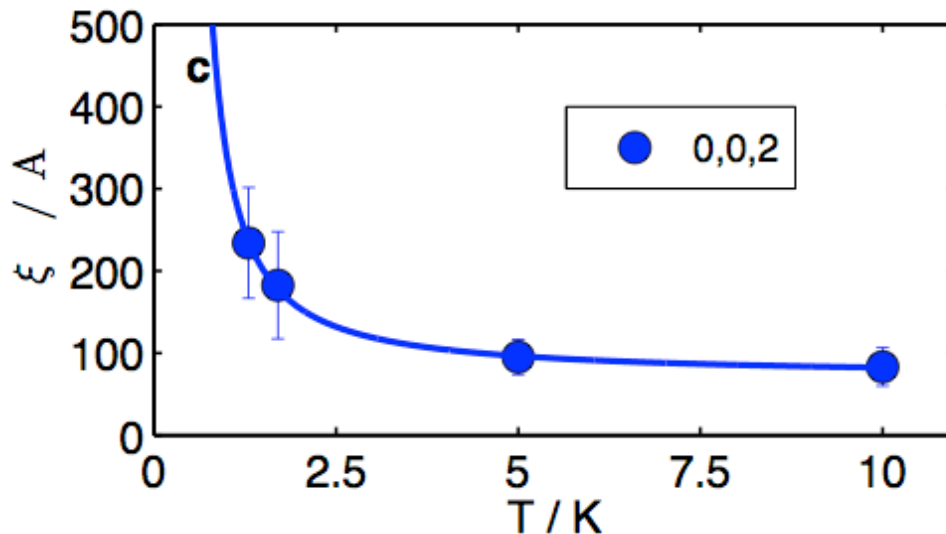
#### Magnetic Coulomb Phase in the Spin Ice $\text{Ho}_2\text{Ti}_2\text{O}_7$

T. Fennell<sup>1,\*</sup>, P. P. Deen<sup>1</sup>, A. R. Wildes<sup>1</sup>, K. Schmalzl<sup>2</sup>, D. Prabhakaran<sup>3</sup>, A. T. Boothroyd<sup>3</sup>, R. J. Aldus<sup>4</sup>,  
D. F. McMorrow<sup>4</sup> and S. T. Bramwell<sup>4</sup>

Published Online 3 September 2009  
*Science* 16 October 2009:  
Vol. 326 no. 5951 pp. 415–417



TRIPLE-AXIS IN12



Again broad region of  $S(Q,\omega)$   
but  
requires variable resolution

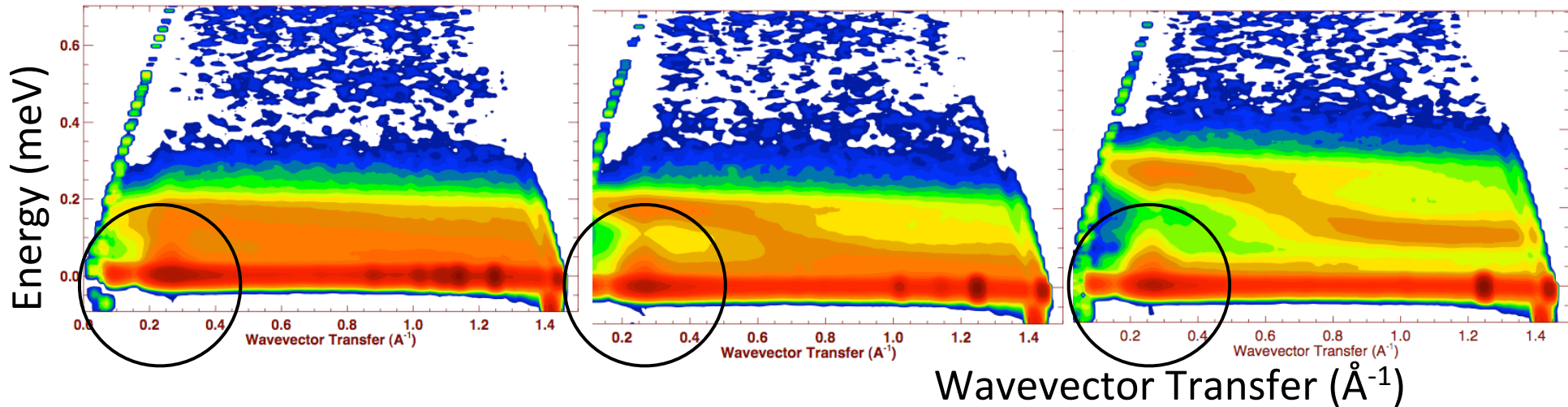
Required the resolution of a triple axis machine to measure the diverging correlation lengths but requires large  $S(Q, \omega)$  for whole image.

- Describes the formation of single and double ice rules defects.  
= magnetic monopoles

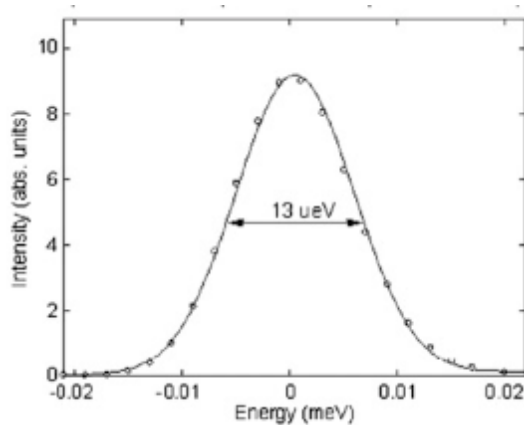


## Extras: Background & resolution : making life easy for the user

(a) Background issues: Cryogenic vacuum - limit Al windows.  
1.0 Tesla                      1.6 Tesla                      2.5 Tesla



(b) Resolution function



Ikeda-Carpenter versus Gaussian, Gaussian please.





- Intense neutron spectra Cold (10 microeV - 30 meV), Thermal (10 - 100 meV) (bi-spectral?)
- Tunable resolution (Gaussian..)
- Excellent PSD coverage including low angles
- Fully integrated (but optional) polarisation analysis (XYZ, CRYOPAD)
- Very low background - high S/N
- Possibility to work with advanced sample environments
- Small samples - focussing optics (Pressure)
- Sample characterisation (in-situ and out)
- Theory groups:  
Nice images are being produced but we are unable to model the systems - theory groups are not a luxury but will be a necessity!